

The Lecture Contains:

 [Thermal Boundary layer equation](#)

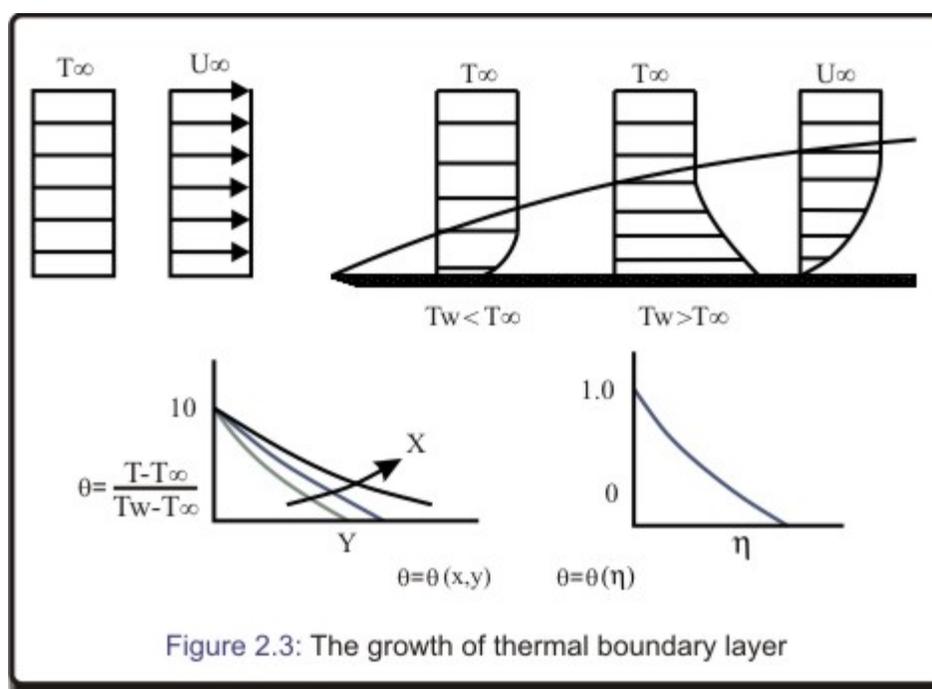
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Thermal Boundary layer equation :-

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \frac{k}{\rho c_p} \frac{\partial^2 T}{\partial y^2} \quad (2.30)$$

$$u = U_\infty f'(\eta), v = -\frac{1}{2} \sqrt{\frac{\nu U_\infty}{x}} \{f(\eta) - \eta f'(\eta)\} \quad (2.31)$$

$$\theta = \frac{T - T_\infty}{T_w - T_\infty} \quad (2.32)$$



describes the growth of thermal boundary layer on a heated at plate. The boundary condition for the situation described above is $\theta(0) = 1$, $\theta(\infty) = 0$. Therefore, $T = T_\infty + \theta(T_w - T_\infty)$.

Let us assume $T_w = T_w(x)$ [a general case]

On substituting $\eta = (y/\sqrt{\frac{\nu x}{U_\infty}})$, we can write different derivatives as,

$$\frac{\partial T}{\partial x} = \theta \frac{dT_w}{dx} + (T_w - T_\infty) \left(-\frac{y}{2x}\right) \sqrt{\frac{U_\infty}{\nu x}} \frac{d\theta}{d\eta} \quad (2.33)$$

$$\frac{\partial T}{\partial y} = \frac{\partial T}{\partial \eta} \frac{\partial \eta}{\partial y} = (T_w - T_\infty) \frac{\partial \theta}{\partial \eta} \sqrt{\frac{U_\infty}{\nu x}} \quad (2.34)$$

$$\frac{\partial^2 T}{\partial y^2} = \frac{\partial}{\partial \eta} \left[\frac{\partial T}{\partial y} \right] \frac{\partial \eta}{\partial y} = (T_w - T_\infty) \frac{\partial^2 \theta}{\partial \eta^2} \frac{U_\infty}{\nu x} \quad (2.35)$$

On substituting the derivatives in thermal boundary layer **equation and multiplying by** $\frac{x}{U_\infty(T_w - T_\infty)}$ **on both sides of the equation**, we get :

$$\frac{f' \theta x}{(T_w - T_\infty)} \frac{dT_w}{dx} - \frac{f' y}{2} \sqrt{\frac{U_\infty}{\nu x}} \theta' - \frac{(f - \eta f')}{2} \theta' = \frac{k}{\rho c_p \nu} \theta'' \quad (2.36)$$

$$f' \theta x \frac{d}{dx} [\ln(T_w - T_\infty)] - \frac{f' \eta \theta'}{2} - \frac{(f - \eta f') \theta'}{2} = \frac{k}{\mu c_p} \theta'' \quad (2.37)$$

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Dividing by $f' \theta$

$$x \frac{d}{dx} [\ln(T_w - T_\infty)] - \frac{f\theta'}{2f'\theta} = \frac{k}{\mu C_p} \theta'' \frac{1}{f'\theta} \quad (2.38)$$

This equation could be put in the form :

$$\frac{x(dT_w/dx)}{T_w(x) - T_\infty} = \frac{k\theta''(\eta)}{\mu c_p f'(\eta)\theta(\eta)} + \frac{f(\eta)\theta'(\eta)}{2f'(\eta)\theta(\eta)} = \lambda \quad (2.39)$$

This equation can be solved by using the method of separation of variables. This leads to :

$$\frac{x(dT_w/dx)}{T_w(x) - T_\infty} = \lambda \quad (2.40)$$

Wall temperature has to follow this relation :

$$d \{ \ln(T_w - T_\infty) \} = \frac{dx}{x} \lambda \quad (2.41)$$

Integrating we get,

$$\ln(T_w - T_\infty) = \lambda \ln x + C \quad (2.42)$$

So,

$$(T_w - T_\infty) = Cx^\lambda$$

When

$$\lambda = 0, T_w = T_\infty + C \quad (\text{constant wall temperature})$$

Again

$$\frac{k\theta''(\eta)}{\mu c_p f'(\eta)\theta(\eta)} + \frac{f(\eta)\theta'(\eta)}{2f'(\eta)\theta(\eta)} = \lambda \quad (2.43)$$

$$\frac{\theta''(\eta)}{Pr} + \frac{1}{2}f\theta' = \lambda f'\theta \quad (2.44)$$

$$\theta'' + \frac{Pr}{2}f\theta' = \lambda Pr f'\theta \quad (2.45)$$

$$\theta'' + \frac{Pr}{2}f\theta' - \lambda Pr f'\theta = 0 \quad (2.46)$$

In this equation f is known from **Blasius solution**. The boundary conditions are:

$$y = 0, T = T_w \rightarrow \eta = 0, \theta = 1$$

$$y = \infty, T = T_\infty \rightarrow \eta = \infty, \theta = 0$$

$$T = T_\infty + \theta(T_w - T_\infty)$$

where

$$(T_w - T_\infty) = Cx^\lambda$$

The final form of the boundary conditions are $\theta(0) = 1$ and $\theta(\infty) = 0$.

Let $Y_1 = \theta$ and $Y_2 = \theta'$ then

$$Y_2' = -\frac{Pr}{2}fY_2 + \lambda Pr f'Y_1 \quad (2.47)$$

with $Y_1(0) = 1$ and $Y_1(\infty) = 0$

The numerical solution using Runge-Kutta method can be obtained via following steps :

1- Guess a value of $Y_2(0)$ or $\theta'(0)$

2- Solve for Y_1, Y_2

3- Check if $Y_1(\infty) = 0$, i.e., $\theta(\infty) = 0$?

4- If yes, stop. The calculated $\theta(\eta)$'s are correct solution

5- If no, correct and $\theta'(0)$ or $Y_2(0)$ repeat the calculation

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Analytical solution, if $\lambda = 0$

$$\theta'' + \frac{Pr}{2} f \theta' = 0 \quad (2.48)$$

$$\frac{\theta''}{\theta'} = -\frac{Pr}{2} f \quad (2.49)$$

or,

$$\ln \theta' = -\frac{Pr}{2} \int f(\eta) d\eta + \ln C$$

or,

$$\theta' = C e^{-\frac{Pr}{2} \int f(\eta) d\eta}$$

or,

$$\theta = C \int e^{-\frac{Pr}{2} \int f(\eta) d\eta} d\eta + D \quad (2.50)$$

or,

$$\theta(\eta) - \theta(0) = C \int_0^\eta e^{-\frac{Pr}{2} \int_0^\eta f(\eta) d\eta} d\eta$$

or,

$$\theta(\eta) = 1 + C \int_0^\beta e^{-\frac{Pr}{2} \int_0^\beta f(r) dr} d\beta$$

Now $\theta(\infty) = 0$ will result in :

$$-1 = C \int_0^{\infty} e^{-\frac{Pr}{2} \int_0^{\beta} f(r) dr} d\beta \quad (2.51)$$

$$C = -\frac{1}{\int_0^{\infty} e^{-\frac{Pr}{2} \int_0^{\beta} f(r) dr} d\beta} \quad (2.52)$$

$$\theta(\eta) = 1 - \frac{\int_0^{\eta} e^{-\frac{Pr}{2} \int_0^{\beta} f(r) dr} d\beta}{\int_0^{\infty} e^{-\frac{Pr}{2} \int_0^{\beta} f(r) dr} d\beta} \quad (2.53)$$

$Nu_x = \text{Nusselt number} = \frac{hx}{k}$

$$h = \frac{-k(dT/dy)_{y=0}}{(T_w - T_{\infty})} = \frac{-k \left[\frac{\partial T}{\partial \eta} \frac{\partial \eta}{\partial y} \right]_{\eta=0}}{(T_w - T_{\infty})} \quad (2.54)$$

We know ,

$$T = T_{\infty} + \theta (T_w - T_{\infty}) \Rightarrow \frac{\partial T}{\partial \eta} = \frac{\partial \theta}{\partial \eta} (T_w - T_{\infty})$$

$$h = -k \frac{\partial \theta}{\partial \eta} \frac{\partial \eta}{\partial y} \Big|_{\eta=0} = -k \theta'(0) \sqrt{\frac{U_{\infty}}{\nu x}} \quad (2.55)$$

$$Nu_x = \frac{hx}{k} = -\theta'(0) \sqrt{\frac{U_{\infty} x}{\nu}} \quad (2.56)$$

$$\frac{Nu_x}{(Re_x)^{1/2}} = -\theta'(0) \simeq 0.332 Pr^{1/3} \quad (0.6 \leq Pr \leq 10) \quad (2.57)$$

The solution was obtained by Pohlhausen .