

Module 3 : Radiation

Lecture 29 : Discrete Ordinates Methodology

Objectives

In this class:

- The solution for the RTE using the Discrete Ordinates Methodology is discussed.
- No analytical solutions will be discussed in this course for the RTE

Discrete Ordinates Methodology (DOM)-1

- The main problem with the RTE is the occurrence of the integral term and differential terms together.
- In the Discrete Ordinates Method the integral term in the RTE is replaced by a summation.

$$\int f(x)dx \cong \sum_{i=1}^n w_i f(x_i) \quad (29.1)$$

- Regular numerical schemes can then be used to solve for the intensities

DOM-2

- Depending on the number of directions chosen the methods are referred to as the S_2 , S_4 , S_6 S_N methods.
- The weights depend on the directions chosen and are computed mathematically. Weights are different for 1D, 2D and 3D situations. Weights and directions ($\mu = \cos\theta$) for 1D and are given in the table on next slide from ref[5].

DOM- 1 D weights & Directions

S_N	Weights	$\pm \mu$
S_2	1	0.50000
S_4	1/2	0.211325
	1/2	0.788675
S_6	1/3	0.146446
	1/3	0.500000
	1/3	0.853554
S_8	1/4	0.102672
	1/4	0.406205
	1/4	0.593795
	1/4	0.897327

- Weights for 1D case. Multiply weights by 2π .
(Taken from Ref.[5])

DOM: Summation terms

- Consider the radiative transport equation:

$$\frac{dl}{ds} = \varepsilon I_b - \alpha I - \rho I + \int_0^{4\pi} \rho I \frac{\hat{\Phi}(\theta_i, \theta_e)}{4\pi} d\theta_i \quad (29.2)$$

- The integral term is represented by:

$$\begin{aligned} \int_0^{4\pi} \rho(s) I(s, \theta_i) \frac{\hat{\Phi}(s, \theta_i, \theta_e)}{4\pi} d\theta_i \\ = \frac{\rho(s)}{4\pi} \sum_{j=1}^n w_j \cdot I(s, \theta_j) \cdot \hat{\Phi}(s, \theta_j, \theta_e) \end{aligned} \quad (29.3)$$

DOM: Directions-1

- The direction 's' is along that specified by the discrete ordinates method. We now look at the cartesian coordinate system. An arbitrary unit vector along a direction which makes angle θ with the z-axis and whose projection on the x-y plane makes an angle ψ with the x axis is given by:

$$\sin \theta \sin \psi \hat{i} + \sin \theta \cos \psi \hat{j} + \cos \theta \hat{k}$$

DOM: Directions-2

- Therefore change in intensity per unit distance is given as:

$$\frac{dI}{ds} = \frac{\partial I}{\partial x} \frac{dx}{ds} + \frac{\partial I}{\partial y} \frac{dy}{ds} + \frac{\partial I}{\partial z} \frac{dz}{ds} \quad (29.4)$$

- And:

$$\frac{dx}{ds} = \sin \theta \sin \psi, \frac{dy}{ds} = \sin \theta \cos \psi, \frac{dz}{ds} = \cos \theta \quad (29.5)$$

DOM: Directions-3

- Therefore:

$$\frac{dI}{ds} = \frac{\partial I}{\partial x} \sin \theta \sin \psi + \frac{\partial I}{\partial y} \sin \theta \cos \psi + \frac{\partial I}{\partial z} \cos \theta \quad (29.6)$$

- Therefore since the discrete ordinate directions are known one can take a cartesian grid and determine the gradients in the x, y and z directions and then compute the change in the intensity in a given direction.

DOM: Directions-4

- For the sake of simplicity define the following:

$$\begin{aligned} \sin \theta \sin \psi &\equiv l \\ \sin \theta \cos \psi &\equiv m \\ \cos \theta &\equiv n \end{aligned} \quad (29.6a)$$

- Therefore:

$$\frac{dI}{ds} = \frac{\partial I}{\partial x} l + \frac{\partial I}{\partial y} m + \frac{\partial I}{\partial z} n \quad (29.6b)$$

DOM: Formulation-1

- Integrate equⁿ (29.2) over a typical control volume:

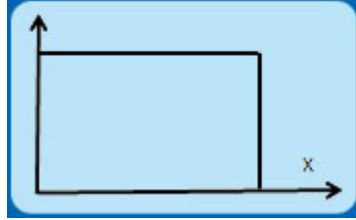
$$\begin{aligned} \iiint \frac{dI}{ds} dV &= \iiint \varepsilon I_b dV - \iiint (\alpha + \rho) I dV \\ &\quad \iiint \left(\int_0^{4\pi} \rho I \frac{\hat{\Phi}(\theta_i, \theta_e)}{4\pi} d\theta_i \right) dV \end{aligned} \quad (29.7)$$

Using the Gauss divergence theorem the first term of the above equation is written as:

$$\iiint \frac{dI}{ds} dV = \iint I l dA_x + \iint I m dA_y + \iint I n dA_z \quad (29.8)$$

DOM: Formulation-2

- Equⁿ (29.8) is substituted in equⁿ (29.7) for the first term.
- The boundary conditions are given and the temperature within the enclosure is to be determined. The entire volume is divided into a number of smaller volumes.



DOM: Formulation-3

- The following nomenclature is used during the integration process:
 A_N , A_S , A_E , A_W : Areas of the North, South, East and West faces respectively.

$I_{P,Q}^k$: the Intensity in the ' k^{th} ' direction at the P, Q location/node

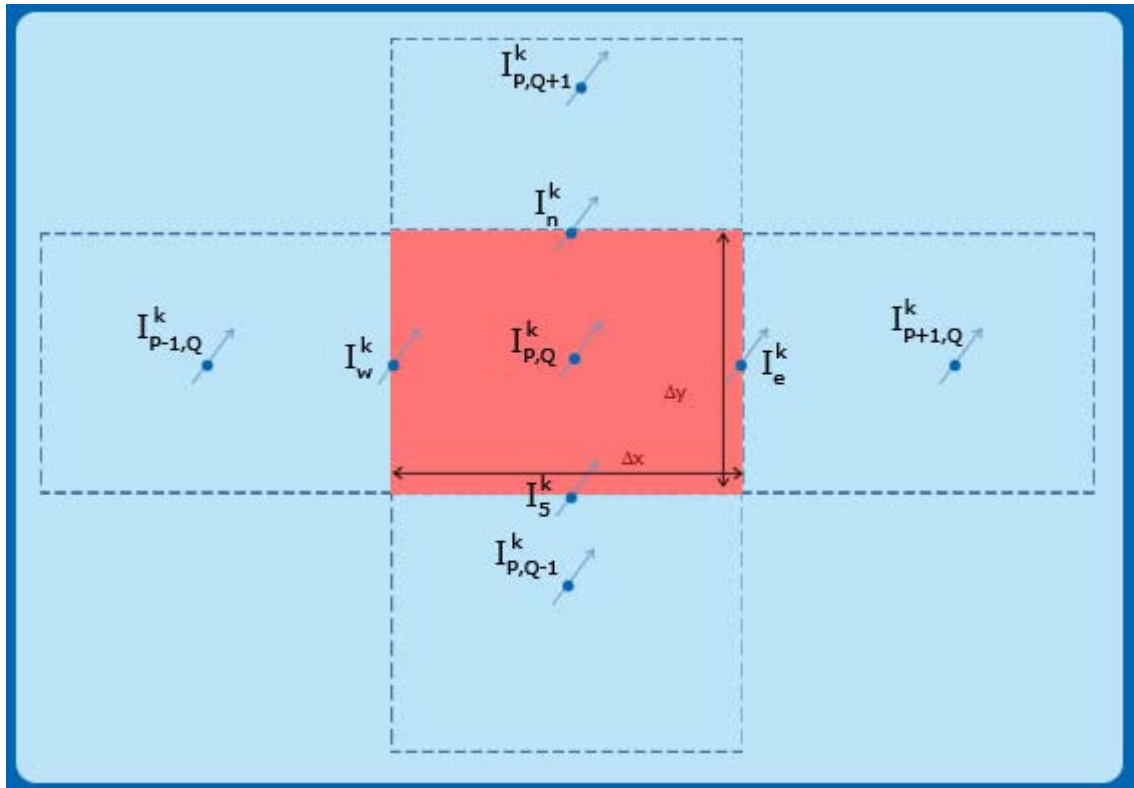
Direction cosines:

$$\frac{dx}{ds} = \sin \theta \sin \psi \equiv l; \frac{dy}{ds} = \sin \theta \cos \psi \equiv m; \frac{dz}{ds} = \cos \theta \equiv n$$

DOM: Formulation-4

- Intensity is different in different directions. Consider the direction k . Consider an arbitrary point **P**, **Q** within the volume. The volume averaged intensity $I_{P,Q}^k$ for the control volume is considered to exist at this point. Intensity at neighbours are $I_{P+1,Q}^k, I_{P-1,Q}^k, I_{P,Q+1}^k, I_{P,Q-1}^k$ and the intensity at the control volume faces are given as $I_e^k, I_w^k, I_n^k, I_s^k$. All these are shown in the figure on the next page where the control volume is in red.

DOM: Formulation-5



DOM: Formulation-6

- For a 2D situation the variation of I^k in the \mathbf{x} direction does not exist, therefore equⁿ (29.8) is rewritten as:

$$\begin{aligned} \iiint \frac{dI^k}{ds} dV &= \iint I^k m dA_y + \iint I^k n dA_z \\ &= m I_e^k A_e - m I_w^k A_w + n I_n^k A_n - n I_s^k A_s \end{aligned} \quad (29.9)$$

- The intensities at the faces are not known and therefore some sort of an interpolation has to be used to relate the cell face and volume intensities.

DOM: Formulation-7

- A common interpolation is:

$$\begin{aligned} I_{P,Q}^k &= \eta I_n^k + (1-\eta) I_s^k ; I_{P,Q}^k = \eta I_e^k + (1-\eta) I_w^k \\ \Rightarrow I_n^k &= \frac{I_{P,Q}^k - (1-\eta) I_s^k}{\eta} ; I_e^k = \frac{I_{P,Q}^k - (1-\eta) I_w^k}{\eta} \end{aligned} \quad (29.10)$$

- Substitute equⁿ (29.10) in equⁿ (29.9) to get:

$$\begin{aligned} \iiint \frac{dI^k}{ds} dV &= m \frac{I_{P,Q}^k - (1-\eta) I_w^k}{\eta} A_e - m I_w^k A_w + n \frac{I_{P,Q}^k - (1-\eta) I_s^k}{\eta} A_n - n I_s^k A_s \\ &= I_{P,Q}^k \left(\frac{n A_n}{\eta} + \frac{m A_e}{\eta} \right) - n I_s^k \left(A_s + A_n \frac{1-\eta}{\eta} \right) - I_w^k m \left(A_w + \frac{1-\eta}{\eta} A_e \right) \end{aligned} \quad (29.11)$$

DOM: Formulation-8

- Substitute equⁿ (29.11) in equⁿ (29.7) to get

$$\begin{aligned} &I_{P,Q}^k \left(\frac{m A_e}{\eta} + \frac{n A_n}{\eta} \right) - n I_s^k \left(A_s + A_n \frac{1-\eta}{\eta} \right) - I_w^k m \left(A_w + \frac{1-\eta}{\eta} A_e \right) \\ &= \varepsilon I_b - (\alpha + \rho) I_{P,Q}^k V_{P,Q} + \frac{\rho(s)}{4\pi} \sum_{j=1}^n w_j I(s, \theta_j) \cdot \hat{\Phi}(s, \theta_j, \theta_e) V_{P,Q} \end{aligned} \quad (29.12)$$

- This equation is now simplified so that the intensity at the point \mathbf{P}, \mathbf{Q} is explicitly obtained.

DOM: Formulation-9

- Simplifying equⁿ (29.12) gives the expression to use for solving for the intensity at the point \mathbf{P}, \mathbf{Q} :

$$I_{P,Q}^k = \frac{\begin{aligned} &I_w^k m (\eta A_w + (1-\eta) A_e) + n I_s^k (\eta A_s + (1-\eta) A_n) + \\ &\eta \varepsilon I_b V_{P,Q} + \eta \frac{\rho(s)}{4\pi} \sum_{j=1}^n w_j I(s, \theta_j) \cdot \hat{\Phi}(s, \theta_j, \theta_e) V_{P,Q} \end{aligned}}{(m A_e + n A_n + \eta (\alpha + \rho) V_{P,Q})} \quad (29.13)$$

- The term marked red is the numerator and that marked blue is the denominator

DOM: Formulation-10

- So far the boundary conditions have not been discussed. At the boundary the intensity will have contributions due to emission and reflection of incident radiation. Recall the equⁿ (26.15) where the definition of intensity was given. This was for intensity of radiation from a surface losing energy. The same magnitude incoming intensity can be visualized as supplying the same energy to a very similar surface.

DOM: Formulation-11

- The incident energy on a surface is therefore:

$$\frac{E_{\text{total}}}{dA_i} = \int I \cos \theta d\omega \quad (29.14)$$

- Assuming the surface to be diffuse, the reflected intensity will be equal in all directions. The reflected intensity is therefore:

$$I = \frac{E_{\text{total}}}{\pi dA_i} = \frac{(1-\varepsilon)}{\pi} \int I \cos \theta d\omega \quad (29.15)$$

DOM: Formulation-12

- The boundary condition at the wall is therefore written as

$$I_{\text{wall}}^k = \varepsilon I_{\text{wall}} + \frac{(1-\varepsilon)}{\pi} \sum_{j=1}^{n/2} \cos(\theta_j) w_j I^j \quad (29.16)$$

- Note that the θ_j is the angle between the normal to the surface and the I^j direction.
- Also note that the summation is only till $n/2$ i.e. only directions towards the surface are considered in the calculations. Only what falls on a surface can get reflected

DOM: Formulation-13

- The steady state radiative energy equation assuming convection, conduction and internal heat generation are negligible is:
- Energy emitted = energy absorbed

$$\text{Energy emitted} = \int \varepsilon I_b d\omega = 4\pi\varepsilon \frac{\sigma T^4}{\pi} \quad (29.17)$$

$$\text{Energy absorbed} = \alpha \int I d\omega \quad (29.18)$$

- Therefore

$$\varepsilon \int I d\omega = 4\varepsilon\sigma T^4 \Rightarrow T^4 = \frac{1}{4\sigma} \int I d\omega \quad (29.19)$$

Recap**In this class:**

- The solution for the RTE using the Discrete Ordinates Methodology is discussed.
- No analytical solutions will be discussed in this course for the RTE