

Module 2 : Convection

Lecture 16 : Fully Developed Pipe flow with Constant Wall temperature and Heat Flux

Objectives

In this class:

- The fully developed temperature profile for uniform circumferential heating with constant flux and the corresponding Nusselt number are derived
- For the constant temperature boundary condition, the methodology is established to the point where numerical calculations are required.

Steady, Incomp., Θ symmetric, pipe Flow- Fully Dev. Temp., Vel. Profile-10

- The equⁿ (15.18) therefore becomes:

$$u^* \frac{\partial T}{\partial z^*} = \frac{kZ}{r_0^2} \frac{1}{\rho c_p \bar{u} r^*} \frac{\partial}{\partial r^*} r^* \frac{\partial T}{\partial r^*} + \frac{k}{Z \rho c_p \bar{u}} \frac{\partial^2 T}{\partial z^{*2}} \quad (16.1)$$

- The above equation suggests the choice for making the 'z' direction non-dimensional. The dimensional combination of Z can be obtained from either the first or the second term on the RHS of equⁿ (16.1).

Steady, Incomp., Θ symmetric, pipe Flow- Fully Dev. Temp., Vel. Profile-11

- Look at the first term to get

$$Z = \frac{\rho c_p \bar{u} r_0^2}{k} = \frac{(\bar{u} r_0) r_0}{\alpha} = \frac{\text{Re Pr } r_0}{2} \quad (16.2)$$
$$\text{Re} = \frac{(\bar{u} 2r_0)}{\nu}$$

- Look at second term to get:

$$Z = \frac{k}{\rho c_p \bar{u}} = \frac{r_0^2 k}{r_0^2 \rho c_p \bar{u}} = \frac{2r_0}{\text{Re Pr}} \quad (16.3)$$

- While 'Z' can be obtained from either (16.2) or (16.3) we will use 'Z' from equⁿ (16.2)

Steady, Incomp., Θ symmetric, pipe Flow- Fully Dev. Temp., Vel. Profile-12

- Now simplify equⁿ (16.1) to get

$$u^* \frac{\partial T}{\partial z^*} = \frac{1}{r^*} \frac{\partial}{\partial r^*} r^* \frac{\partial T}{\partial r^*} + \frac{k}{\rho c_p \bar{u} r_0^2} \frac{\partial^2 T}{\partial z^{*2}} \frac{1}{\rho c_p \bar{u}}$$
$$u^* \frac{\partial T}{\partial z^*} = \frac{1}{r^*} \frac{\partial}{\partial r^*} r^* \frac{\partial T}{\partial r^*} + \frac{4}{(\text{Re Pr})^2} \frac{\partial^2 T}{\partial z^{*2}} \quad (16.4)$$

- Define $\text{RePr} = \text{Pe}$. Pe = Peclet number.
- High Pe means axial conduction is ignored wrt to convection (or radial conduction – since convection term arises from radial conduction from the wall).

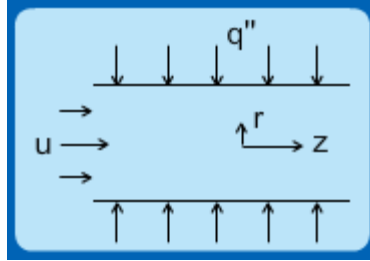
Steady, Incomp., Θ symmetric, pipe Flow- Fully Dev. Temp., Vel. Profile-13

- Peclet number high, implies ignoring the axial conduction term with respect to the radial conduction (or the convection term since radial conduction and convection terms are of similar magnitude)
- Now let us make an additional assumption that the Peclet number is high and so ignore the axial conduction term. The final equation is therefore:

$$\rho c_p u \frac{\partial T}{\partial z} = \frac{k}{r} \frac{\partial}{\partial r} \left(r \frac{\partial T}{\partial r} \right) \quad (16.5)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Flux-1

- Solve equⁿ (16.5) with appropriate boundary conditions
- Consider flow through a pipe with uniform wall flux at the wall. Assume flow is fully developed.



Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Flux-2

- Use the nondimensional temperature profile and thermally fully developed flow:

$$\theta(r) = \frac{T(z, r) - T_w(z)}{T_b(z) - T_w(z)}$$

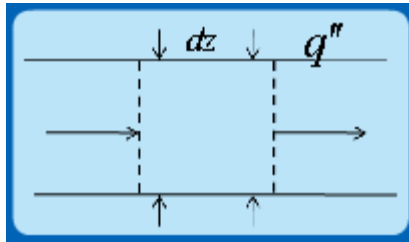
$$T = T_w + \theta(T_b - T_w)$$

- Equⁿ (16.5) needs the temperature gradient in the z direction. Therefore:

$$\frac{\partial T}{\partial z} = \frac{dT_w}{dz} + \theta(r) \left(\frac{dT_b}{dz} - \frac{dT_w}{dz} \right) \quad (16.6)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Flux-3

- Consider a small slice of the pipe and use the First Law of Thermodynamics for an open system



$$q''(2\pi r_0 dz) = \dot{m} c_p dT_b$$

$$\frac{dT_b}{dz} = \frac{2q''}{\rho \bar{u} r_0 c_p} \quad (16.7)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Flux-4

- For fully developed thermal profile heat transfer coefficient is constant. Wall heat flux is also constant. Therefore:

$$q_w = h(T_w - T_b)$$

$$\Rightarrow (T_w - T_b) = \text{constant} \Rightarrow \frac{dT_w}{dz} = \frac{dT_b}{dz} \quad (16.8)$$

- Substitute equⁿ (16.8) in equⁿ (16.6) gives:

$$\frac{\partial T}{\partial z} = \frac{dT_w}{dz} \quad (16.9)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Flux-5

- The governing equⁿ (16.5) therefore becomes:

$$\rho c_p u \frac{dT_b}{dz} = \frac{k}{r} \frac{\partial}{\partial r} \left(r \frac{\partial T}{\partial r} \right) \quad (16.10)$$

- Substitute equⁿ (16.7) in equⁿ (16.10) to get:

$$\begin{aligned} \rho c_p u \frac{2q''}{\rho \bar{u} r_0 c_p} &= \frac{k}{r} \frac{\partial}{\partial r} r \frac{\partial T}{\partial r} \\ \Rightarrow \frac{u}{\bar{u}} \frac{2q''}{r_0} &= \frac{k}{r} \frac{\partial}{\partial r} r \frac{\partial T}{\partial r} \end{aligned} \quad (16.11)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Flux-6

- Use the following non-dimensional variables:

$$r^* = \frac{r}{r_0}, u^* = \frac{u}{\bar{u}} = 2(1 - r^{*2}), \theta = \frac{T - T_w}{T_b - T_w} \quad (16.12)$$

- Non-dimensional energy equation with the boundary conditions is thus:

$$4(1 - r^{*2}) \frac{q''}{k r_0} = \frac{(T_b - T_w)}{r_0^2 r^*} \frac{d}{dr^*} r^* \frac{d\theta}{dr^*} \quad (16.13)$$

$$r^* = 0, \frac{\partial T}{\partial r} = 0 \Rightarrow \left. \frac{\partial \theta}{\partial r^*} \right|_{r^*=0} = 0 \quad (16.14)$$

$$r = r_0, T = T_w \Rightarrow \theta|_{r^*=1} = 0 \quad (16.15)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Flux-7

- Integrating equⁿ (16.13) gives

$$\begin{aligned} \int 4q'' \frac{r_0}{k} (1 - r^{*2}) r^* dr^* &= (T_b - T_w) \int \frac{d}{dr^*} r^* \frac{d\theta}{dr^*} \\ q'' \frac{r_0}{k} \left(\frac{r^{*2}}{2} - \frac{r^{*4}}{4} \right) + C_1 &= (T_b - T_w) r^* \frac{d\theta}{dr^*} \end{aligned} \quad (16.16)$$

- Use condition in equⁿ (16.14) to get $C_1 = 0$. Now integrate equⁿ (16.16) to get:

$$T - T_w = \int 4q'' \frac{r_0}{k} \left(\frac{r^{*2}}{2} - \frac{r^{*4}}{4} \right) dr^* \Rightarrow T - T_w = \frac{q'' r_0}{k} \left[r^{*2} - \frac{r^{*4}}{4} \right] + C_2$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Flux-7

- Use condition in equⁿ (16.15) to get:

$$T - T_w = \frac{q'' r_0}{k} \left[r^{*2} - \frac{r^{*4}}{4} - \frac{3}{4} \right] \quad (16.17)$$

- The temperature profile is known and thus the bulk temperature can be evaluated using:

$$T_b = \frac{\int \dot{m} c_p T}{\int \dot{m} c_p} = \frac{\int 2\pi r dr u T}{\pi r_0^2 \bar{u}} \quad (16.18)$$

$$\Rightarrow T_b - T_w = \frac{\int 2\pi r dr u T}{\pi r_0^2 \bar{u}} - T_w \quad (16.19)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Flux-8

- Equⁿ (16.19) can be written as:

$$T_b - T_w = \frac{\int 2\pi r dr u (T - T_w)}{\pi r_0^2 \bar{u}} \quad (16.20)$$

- T_w is taken inside integral on the RHS of equⁿ (16.19) since it is constant at any 'z' location
- Use equⁿ (16.17) in equⁿ (16.20) and integrate:

$$\begin{aligned} \therefore T_b - T_w &= \int 4(1-r^{*2}) r^{*} \frac{q'' r_0}{k} \left(r^{*2} - \frac{r^{*4}}{4} - \frac{3}{4} \right) dr^{*} \\ &= 4 \frac{q'' r_0}{k} \left[\frac{1}{4} - \frac{1}{24} - \frac{3}{8} - \frac{1}{6} + \frac{1}{32} - \frac{3}{16} \right] = -\frac{11}{24} q'' \frac{r_0}{k} \end{aligned} \quad (16.21)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Flux-9

- Substitute (16.21) in (16.17) to get

$$T - T_b = \frac{q'' r_0}{k} \left[r^{*2} - \frac{r^{*4}}{4} - \frac{7}{24} \right] \quad (16.22)$$

- Bulk temperature at any 'z' location can be written in terms of the inlet temperature using equⁿ (16.7) and therefore the above equation can also be written as:

$$T = T_i + \frac{2q'' z}{\rho \bar{u} r_0 c_p} + \frac{q'' r_0}{k} \left[r^{*2} - \frac{r^{*4}}{4} - \frac{7}{24} \right] \quad (16.23)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Flux-10

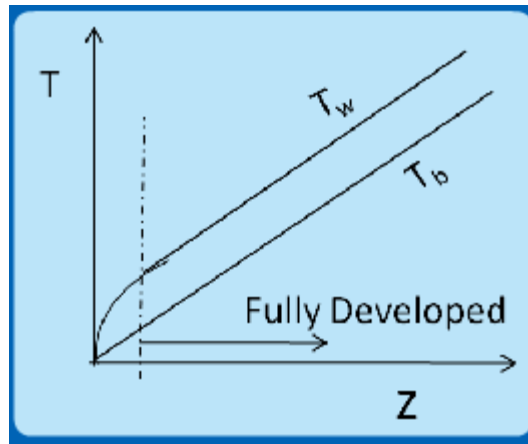
- Now evaluate the heat transfer coefficient using equⁿ (16.21):

$$\begin{aligned} -q'' &= h(T_b - T_w) = h \left(-\frac{11}{24} q'' \frac{r_0}{k} \right) = h \frac{11}{48} q'' \frac{2r_0}{k} \\ Nu &= \frac{h(2r_0)}{k} = \frac{48}{11} = 4.36 \end{aligned}$$

- For constant wall flux case the Nusselt number is a constant.

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Flux-11

- Qualitative variation is shown as a function of axial location
- Note wall and bulk temperatures are parallel in the Fully Developed Zone (equⁿ (16.8))



Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Temp-1

- Slightly more difficult since $\frac{\partial T}{\partial z}$ cannot be replaced by a constant. Governing equation remains the same

$$\rho c_p u \frac{\partial T}{\partial z} = \frac{k}{r} \frac{\partial}{\partial r} r \frac{\partial T}{\partial r} \quad (16.5)$$

- Try to see if the non-dimensional temperature gives some information

$$\begin{aligned} \frac{T - T_w}{T_b - T_w} = \theta(r) &\Rightarrow \frac{\partial T}{\partial z} = \theta(r) \frac{dT_b}{dz} \\ &\Rightarrow \frac{\partial T}{\partial r} = (T_b - T_w) \frac{d\theta}{dr} \end{aligned} \quad (16.24)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Temp-2

- Using the same non-dimensional variables get the equⁿ (16.4) in which 'Pe' term neglected:

$$u^* \frac{\partial T}{\partial z^*} = \frac{1}{r^*} \frac{\partial}{\partial r^*} r^* \frac{\partial T}{\partial r^*}$$

- Use equⁿ (15. 13), (16.24) in (16.5)

$$2(1 - r^{*2})\theta(r) \frac{dT_b}{dz^*} = \frac{1}{r^*} \frac{\partial}{\partial r^*} \left(r^* \frac{d\theta}{dr^*} \right) (T_b - T_w) \quad (16.25)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Temp-3

- Collect like terms to write

$$\frac{dT_b/dz^*}{(T_b - T_w)} = \frac{\frac{d}{dr^*} r^* \frac{\partial \theta}{\partial r^*}}{2(1 - r^{*2})r^* \theta} = -\lambda \quad (16.26)$$

- In the above equation first term is a function of 'z' only while the second term is a function of 'r' only and therefore equal to a constant. The negative value is to keep the T_b bounded.

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Temp-4

- Since wall temperature is constant, Rewrite the first term of equⁿ (16.26) as:

$$\begin{aligned}\frac{dT_b}{dz} &= \frac{d(T_b - T_w)}{dz} = -\lambda \\ \Rightarrow T_b - T_w &= (T_i - T_w)e^{-\lambda z}\end{aligned}\quad (16.26a)$$

- T_i is the temperature with which fluid enters the pipe at $z = 0$.

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Temp-5

- The second term in the equⁿ (16.26) in the 'r' direction gives:

$$\frac{d}{dr} \left(r \frac{d\theta}{dr} \right) + 2\lambda r (1 - r^2) \theta = 0 \quad (16.27)$$

- The boundary conditions are

$$T = T_w \Big|_{r=1} \Rightarrow \theta = 0 \quad (16.28)$$

$$\frac{\partial T}{\partial r} \Big|_{r=0} = 0 \Rightarrow \theta'(0) = 0 \quad (16.29)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Temp-6

- Another condition is needed since λ is still unknown. This comes from the definition of the bulk temperature:

$$T_b = \frac{\int \rho c_p u T 2\pi r dr}{\pi r_0^2 \bar{u} \rho c_p} \quad (16.30)$$

$$\begin{aligned}&= 2 \int_0^{r_0} \frac{u}{\bar{u}} \left(\frac{r}{r_0} \right) T d \left(\frac{r}{r_0} \right) \\ (T_b - T_w) &= 4 \int_0^1 (1 - r^2) r^* (T - T_w) dr^* \quad (16.31)\end{aligned}$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Temp-7

- Note that T_w is subtracted from both sides in equⁿ (16.31). However, it is taken into the integral sign on the RHS since the term within the integral is unity for constant 'T' and T_w is constant. Term inside the integral being unity for constant 'T' is easier understood by viewing equⁿ (16.30)
- Equⁿ (16.31) is rewritten as:

$$\frac{1}{4} = \int_0^1 (1 - r^2) \theta r^* dr^* \quad (16.32)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Temp-8

- Equations (16.27), (16.28), (16.29), (16.32) are solved together for θ and λ . Analytical solution is difficult and therefore a numerical methodology is adopted.
- A methodology could be guess λ and solve equⁿs (16.27), (16.28), (16.29) for θ . However, the equation (16.32) needs to be satisfied also and since it does not have λ explicitly, using this to get a better guess for λ is not possible. One way is define a new variable $F = \lambda \theta$.

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Temp-9

- The equations (16.27), (16.28), (16.29), (16.32) are therefore rewritten as:

$$\frac{d}{dr^*} \left(r^* \frac{dF}{dr^*} \right) + 2\lambda r^* (1 - r^{*2}) F = 0 \quad (16.33)$$

$$F(0) = 0; F'(1) = 0 \quad (16.34)$$

$$\frac{\lambda}{4} = \int_0^1 (1 - r^{*2}) F r^* dr^* \quad (16.35)$$

- Guess a value for λ , solve for F using equⁿs (16.33), (16.34). Use equⁿ (16.35) to get new λ , solve for F again and continue till convergence, i.e. change in λ is very small.

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Temp-10

- Another methodology:
- Guess two values of λ and solve for θ using equⁿs (16.27), (16.28) and (16.29)
- Use equation (16.32) to compute the integral for two values of λ .
- Use secant method till integral becomes 0.25 to within a specified tolerance
- λ will converge to 3.66

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Temp-11

- Sometimes bulk temperature for constant wall temperature case is calculated for the case when flow is fully developed.
- Energy balance gives:

$$\dot{m} C_p dT_b = h P dz (T_w - T_b) \quad (16.36)$$

- Define: $\Theta = T_w - T_b$
- Equⁿ (16.36) transforms to

$$-\frac{d\Theta}{dz} = \frac{hP\Theta}{\dot{m}C_p} \quad (16.37)$$

Pipe Flow- Fully Developed Temp. and Velocity Profile Constant Wall Temp-12

- Solution for equⁿ (16.37) is:

$$\Theta = \Theta_i e^{-\frac{hPz}{\dot{m}C_p}} \quad (16.38)$$

- The temperature profile is therefore exponential as observed earlier also.
- Bulk temperature variation for the constant wall flux and constant wall temperature can be evaluated by just doing an energy balance.

Recap

In this class:

- The fully developed temperature profile for uniform circumferential heating with constant flux and the corresponding Nusselt number are derived
- For the constant temperature boundary condition, the methodology is established to the point where numerical calculations are required.