

## Module 9: Stiff-Initial Value Systems

## Lecture 33: First order linear systems with constant coefficient

## The Lecture Contains:

We briefly discuss linear systems with constant coefficients and illustrate with an example, how to find the general solution of such a system. This is followed by a discussion on solving these initial value systems, say, by an Euler's method which finally leads to the notion of stiffness.

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The first order system  $\mathbf{y}' = \mathbf{f}(t, \mathbf{y})$ , where  $\mathbf{y}$  and  $\mathbf{f}$  are  $m$ -dimensional vectors, is said to be linear if  $\mathbf{f}(t, \mathbf{y}) = \mathbf{A}(t)\mathbf{y} + \boldsymbol{\phi}(t)$ , where  $\mathbf{A}(t)$  is an  $m \times m$  matrix and  $\boldsymbol{\phi}(t)$  an  $m$ -dimensional vector; if in addition,  $\mathbf{A}(t) = \mathbf{A}$ , a constant matrix, the system is said to be linear with constant coefficients. We require the general solution of such a system

$$\mathbf{y}' = \mathbf{A}\mathbf{y} + \boldsymbol{\phi}(t) \quad (9.1)$$

Let  $\hat{\mathbf{y}}(\mathbf{x})$  be the general solution of the corresponding homogeneous system

$$\mathbf{y}' = \mathbf{A}\mathbf{y} \quad (9.2)$$

If  $\boldsymbol{\psi}(t)$  is any particular solution of (9.1), then  $\mathbf{y}(t) = \hat{\mathbf{y}}(t) + \boldsymbol{\psi}(t)$  is the general solution of (9.1).

A set of  $m$  linearly independent solutions  $\hat{\mathbf{y}}_j(t), j = 1, 2, \dots, m$  of (9.2), is said to form a fundamental system of (9.2), and the most general solution of (9.2) may be written as a linear combination of the members of the fundamental system. It is easily seen that  $\hat{\mathbf{y}}_j(t) = e^{\lambda_j t} \mathbf{C}_j$ , where  $\mathbf{C}_j$  is an  $m$ -dimensional vector, is a solution of (9.2) if  $\lambda_j \mathbf{C}_j = \mathbf{A} \mathbf{C}_j$ , that is, if  $\lambda_j$  is an eigen value of  $\mathbf{A}$  and  $\mathbf{C}_j$  is the corresponding eigen vector. We consider only the case where  $\mathbf{A}$  possesses  $m$  distinct possibly complex, eigen values  $\lambda_j, j = 1, 2, \dots, m$ . The corresponding eigen vectors  $\mathbf{C}_j, j = 1, 2, \dots, m$  are then linearly, independent, and it follows that the solutions  $\hat{\mathbf{y}}_j(t) = e^{\lambda_j t} \mathbf{C}_j, j = 1, 2, \dots, m$  form a fundamental system of (9.2). The most general solution of (9.1) is then

$$\mathbf{y}(t) = \sum_{j=1}^m K_j e^{\lambda_j t} \mathbf{C}_j + \boldsymbol{\psi}(t)$$

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Consider the example given by

$$y' = A y; y(0) = [1, 0, -1]^T$$

$$\text{where } A = \begin{bmatrix} -21 & 19 & -20 \\ 19 & -21 & 20 \\ 40 & -40 & -40 \end{bmatrix}$$

The eigen values of  $A$  are the roots of the equation  $\det(A - \lambda I) = 0$  and are found to be  $\lambda_1 = -2, \lambda_2 = -40 + 40i, \lambda_3 = -40 - 40i$ , and are distinct. The corresponding eigen vectors are

$C_1 = [1, 1, 0]^T, C_2 = [1, -1, -2i]^T$  and  $C_3 = [1, -1, 2i]^T$ . The general solution of  $y' = A y$  is

$$y(t) = K_1 + e^{\lambda_1 t} C_1 + K_2 e^{\lambda_2 t} C_2 + K_3 e^{\lambda_3 t} C_3.$$

For this problem,  $\psi(t)$  is identically zero, and the given initial vector  $[1, 0, -1]^T$  can be expressed as the following linear combination of  $C_1, C_2$  and  $C_3$ :

$$[1, 0, -1]^T = \frac{1}{2}[1, 1, 0]^T + \frac{1}{4}(1 - i)[1, -1, -2i]^T + \frac{1}{4}(1 + i)[1, -1, 2i]^T.$$

We thus choose  $K_1 = \frac{1}{2}, K_2 = \frac{1}{4}(1 - i)$  and  $K_3 = \frac{1}{4}(1 + i)$ , giving the solution

$$y(t) = \frac{1}{2} e^{-2t} \begin{bmatrix} 1 \\ 1 \\ 0 \end{bmatrix} + \frac{1}{4}(1 - i) e^{(-40 + 40i)t} \begin{bmatrix} 1 \\ -1 \\ -2i \end{bmatrix} + \frac{1}{4}(1 + i) e^{(-40 - 40i)t} \begin{bmatrix} 1 \\ -1 \\ 2i \end{bmatrix}$$

or

$$u(t) = \frac{1}{2} e^{-2t} + \frac{1}{2} e^{-40t} (\cos 40t + \sin 40t)$$

$$v(t) = \frac{1}{2} e^{-2t} - \frac{1}{2} e^{-40t} (\cos 40t + \sin 40t)$$

$$w(t) = -e^{-40t} (\cos 40t - \sin 40t)$$

where  $y(t) = [u(t), v(t), w(t)]^T$ .

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If we now attempt to solve this problem by Euler's method with  $h = 0.04$  in the range  $0.1 \leq t \leq 1.0$  with  $y(0.1)$  given by the exact solution. We find that for the given problem in the range  $0.1 \leq t \leq 1.0$ , the choice of  $h = 0.04$  causes  $\bar{h}$  to lie outside the region of absolute stability, which is the circle  $\mathbf{R}$  with center  $-1$ , radius 1, and it follows that for  $\bar{h}(= h \lambda)$  to lie within  $\mathbf{R}$  for all three values of  $\lambda$ , we must satisfy  $h < 0.025$ . Note that the eigen values responsible for this severe restriction in  $h$  are  $-40 \pm 40i$ ; that is, the very eigen values whose contributions to the theoretical solution are negligible in the range  $0.1 \leq t \leq 1.0$

On the other hand, consider the IVP

$$y' = Ay \quad y(0.1) = \left[ \frac{1}{2} e^{-0.2}, \frac{1}{2} e^{-0.2}, 0 \right]^T$$

$$\text{where } A = \begin{bmatrix} -2 & 0 & 0 \\ 0 & -2 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$$\text{whose theoretical solution } y(t) = \left[ \frac{1}{2} e^{-2t}, \frac{1}{2} e^{-2t}, 0 \right]^T$$

is, in the range  $0.1 \leq t \leq 1.0$ , virtually indistinguishable from that of the previous problem, is integrated perfectly satisfactorily by Euler's rule with step length  $0.04$ . The Eigen values of the system for this problem are  $-2, -2, 0$  and for absolute stability we require only  $h < 1.0$ .

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We know that the  $m \times m$  linear systems

$$y' = A y + \phi(t)$$

where the matrix  $A$  has distinct eigen values  $\lambda_j$  and corresponding eigen vector  $C_j, j = 1, 2, \dots, m$  has a general solution of the form

$$y(t) = \sum_{j=1}^m K_j e^{\lambda_j t} C_j + \psi(t)$$

Let us assume that  $\text{Re } \lambda_j < 0, j = 1, 2, \dots, m$  then the term  $\sum_{j=1}^m K_j e^{\lambda_j t} C_j \rightarrow 0$  as  $t \rightarrow \infty$  we therefore call this term the transient solution, and call the remaining term  $\psi(t)$  the steady state solution. Let  $\lambda_\mu$  and  $\lambda_\nu$  be two eigen values of  $A$  such that

$$|\text{Re } \lambda_\mu| \geq |\text{Re } \lambda_j| \geq |\text{Re } \lambda_\nu|, j = 1, 2, \dots, m$$

If our aim is to find numerically the steady state solution  $\psi(t)$ , then we must pursue the numerical solution until the slowest decaying exponential in the transient solution, namely  $e^{\lambda_\nu t}$  is negligible. Thus, the smaller  $|\text{Re } \lambda_\nu|$ , the longer will be the range of integration. On the other hand, the presence of eigen values of  $A$  far out to the left in the complex plane will force us to use excessively small step lengths in order that  $\bar{h}$  will lie within the range of absolute stability of the method. The further out such eigen values lie, the more severe is the restriction on step length. A rough measure of this difficulty is the magnitude of  $|\text{Re } \lambda_\mu|$ . If  $|\text{Re } \lambda_\mu| \gg |\text{Re } \lambda_\nu|$ , we are forced into the highly undesirable computational situation of having to integrate numerically over a long range, using a step length which is everywhere excessively small relative to the interval, this is the problem of stiffness. We can make the following somewhat heuristic definition.

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