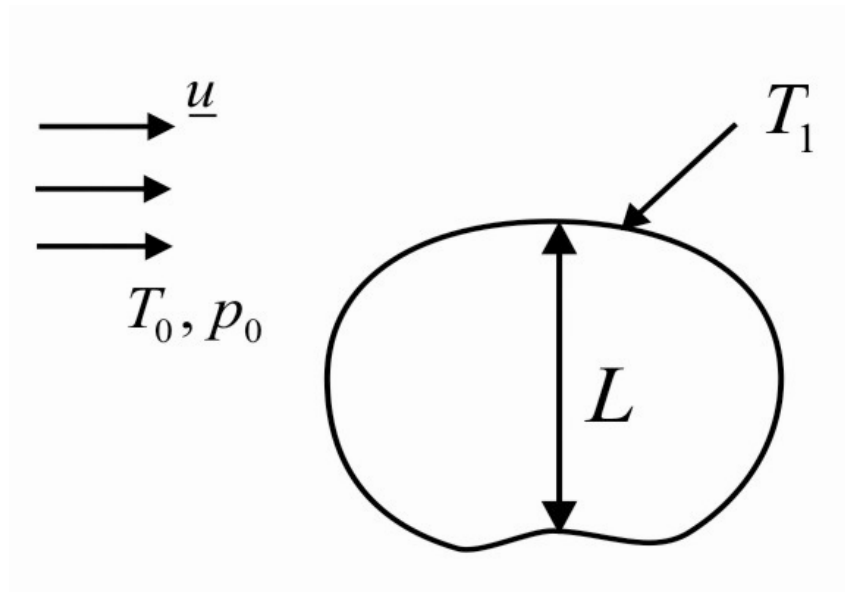


Module 4: Multicomponent Transport

Lecture 37: Dimensional Analysis



Equation of Continuity:

$$\nabla \cdot \underline{u} = \begin{cases} 0 & \text{(incompressible)} \\ -\frac{1}{\rho} \frac{D\rho}{DT} & \text{(compressible)} \end{cases}$$

Navier–Stokes Equation: (Constant physical properties, Newtonian fluid, Boussinesq approximation)

$$\rho \frac{D\underline{u}}{Dt} = \mu \nabla^2 \underline{u} + \begin{cases} -\nabla p + \rho \underline{g} & \text{(forced convection)} \\ -\rho \beta \underline{g} (T - T_0) & \text{(free convection)} \end{cases}$$

Energy Balance Equation:

$$\rho C_p \frac{DT}{Dt} = k \nabla^2 T + \mu \Phi_u + \beta T \frac{Dp}{Dt}$$

Note: $\beta = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$: coefficient of volume expansion

For ideal gases $\Rightarrow \beta = \frac{1}{T}$

Define dimensionless variables (for forced convection)

$$\underline{x}^* = \underline{x}/L, \quad \underline{U}^* = \underline{u}/u, \quad T^* = \frac{T - T_0}{T_1 - T_0}, \quad t^* = \frac{tu}{L}$$

$$P^* = \frac{p - p_0}{\rho u^2}$$

Dimensionless Conservation Equations (Incompressible fluid, Forced Convection case):

$$\nabla^* \cdot \underline{U}^* = 0$$

$$\frac{D\underline{U}^*}{Dt^*} = -\nabla^* P^* + \frac{1}{Fr} \frac{\underline{g}}{g} + \frac{1}{Re} \nabla^{*2} \underline{U}^*$$

$$\frac{DT^*}{Dt^*} = \frac{1}{\text{Re Pr}} \nabla^{*2} T^* + \frac{Br}{\text{Re Pr}} \Phi_U^* + \frac{Br}{\text{Pr}} [\beta T_0 + \beta \Delta T T^*] \frac{DP^*}{Dt^*}$$

Dimensionless Numbers:

$$\text{Re} = \text{Reynolds \#} = \frac{uL}{\nu} = \frac{\rho u^2 / L}{\mu u / L^2} = \frac{\text{Inertial forces}}{\text{Viscous forces}}$$

$$\text{Fr} = \text{Froude \#} = \frac{u^2}{L_g} = \frac{\rho u^2 / L}{\rho g} = \frac{\text{Inertial forces}}{\text{Gravity forces}}$$

$$\text{Re Pr} = \frac{\rho C_p u L}{k} = \text{Peclet(heat)} = \text{Pe}_h = \frac{\ell C_p u (T_0 - T_1) / L}{k (T_0 - T_1) / L^2} = \frac{\text{heat transport by convection}}{\text{heat transport by conduction}}$$

$$\text{Br} = \text{Brinkman \#} = \frac{\mu u^2}{k (T_0 - T_1)} = \frac{\mu (u/L)^2}{k (T_0 - T_1) / L^2} = \frac{\text{heat production by viscous dissipation}}{\text{heat transport by conduction}}$$

(important for turbo-machinery)

$$\text{Ec} = \text{Eckert \#} = \frac{Br}{\text{Pr}} = \frac{\rho u^2}{\rho C_p (T_0 - T_1)} = \frac{\text{Kinetic energy}}{\text{Internal energy}}$$

(important for hypersonic flows)

For free convection there is no readily available reference velocity. So it is more convenient to define

$$\underline{U}^{**} = \frac{uLe}{\mu} = \text{dimensionless velocity}$$

$$t^{**} = \frac{t\mu}{\rho L^2} = \text{dimensionless time}$$

The conservation equation then become

$$\nabla^* \cdot \underline{U}^{**} = 0$$

$$\frac{D\underline{U}^{**}}{Dt^{**}} = \nabla^{*2} \underline{U}^{**} - T^* Gr \frac{\underline{g}}{g}$$

$$\frac{DT^*}{Dt^{**}} = \frac{1}{Pr} \nabla^{*2} T^* + \text{negligible terms}$$

Thus, for free convection only Gr and Pr appear

The Grashof # $Gr = \frac{g\rho^2\beta(T_0 - T_1)L^3}{\mu^2}$

$$\rightarrow \frac{Gr}{Re^2} = \frac{\beta g (T_0 - T_1) L}{u^2} = \frac{\rho \beta g (T_0 - T_1)}{\rho u^2 / L} = \frac{\text{Buoyancy forces}}{\text{Inertial forces}}$$

The above number decides if natural convection is important as compared to forced convection.

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Temperature Distribution: The temperature distribution will be a function of numerous parameters
(say, for forced convection)

$$T^* = F(t^*, \underline{X}^*; \text{Re}, \text{Pr}, Fr, Br, \beta T_0, \beta \Delta T)$$

Usually not important

Not important if there is no free surface flow

So, for a wide range of problems, $T^* = F(t^*, \underline{X}^*, \text{Re}, \text{Pr})$

$$\text{Nusselt number} = \frac{hL}{k} = \frac{Q}{A\Delta T} \frac{L}{k},$$

Where h is the average heat transfer coefficient

$$Q = \int_s -k \frac{\partial T}{\partial \underline{x}} \cdot \underline{n} dA = -kL(T_1 - T_0) \int_{s^*} \frac{\partial T^*}{\partial \underline{x}^*} \cdot \underline{u} dA^*$$

$$= f(\text{geometry}, t^*, \text{Re}, \text{Pr})$$

Analogies between Heat and Mass Transfer: Recall correlations for boundary layer heat transfer and also for heat transfer inside tubes.

$$\text{Pr} \rightarrow \text{Sc} = \frac{\nu}{D}$$

$$\Rightarrow \text{For mass transfer} \quad \text{Pe}_h \rightarrow \text{Pe}_m = \frac{uL}{D}$$

$$\text{Nu} \rightarrow \text{Sh} = \frac{k_m L}{D}$$

$$\text{Gr} \rightarrow \text{Gr}_{AB} = \frac{\rho^2 \mathfrak{S} g (X_{A1} - X_{A0}) L^3}{\mu^2},$$

$$\text{where } \mathfrak{S} = \frac{1}{\rho} \left(\frac{\partial \rho}{\partial x_A} \right)_{p,T}.$$