

## **Module 5:**

# **Mass transfer in turbulent flows**

## **Lecture 39:**

# **Transport in Turbulent Flow: Universal**

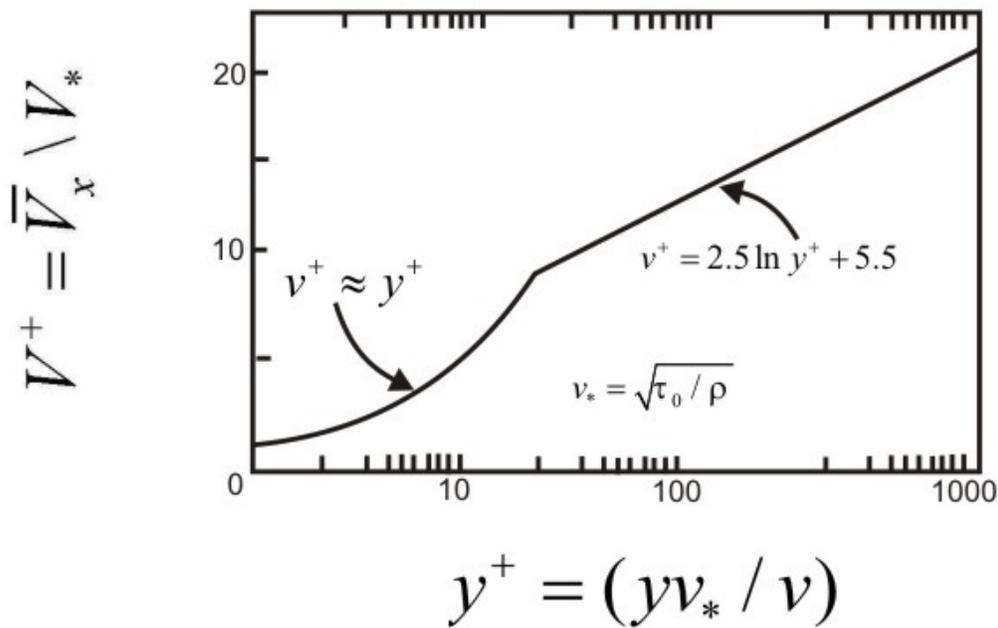
# **Velocity Profile**

## Universal Velocity Profile

For fully developed turbulent flow near a smooth wall: Good for pipe flow as well as boundary layers.

$$v^+ = \frac{\bar{U}_x}{U_*}, y^+ = \frac{yU_*}{\nu} \quad \text{with} \quad u_* = \sqrt{\frac{\tau_0}{\rho}} \quad (4.12)$$

Where  $y$  is the distance from wall and Shear stress  $\tau_0$  deduced from pressure drop measurements.



Near the wall =>  $u^+ \approx y^+ \quad y^+ < 5 \quad (4.13)$

Far from wall =>  $u^+ \cong 2.5 \ln y^+ + 5.5 \quad y^+ > 20 \quad (4.14)$

In the logarithmic region

$$\bar{u}_x = 2.5 \sqrt{\frac{\tau_0}{\rho}} \ln y + \left(2.5 \ln \frac{\sqrt{\tau_0 / \rho}}{\nu} + 5.5\right) \sqrt{\frac{\tau_0}{\rho}} \quad (4.15)$$

Viscosity enters only in the additive constant term

### Eddy viscosity

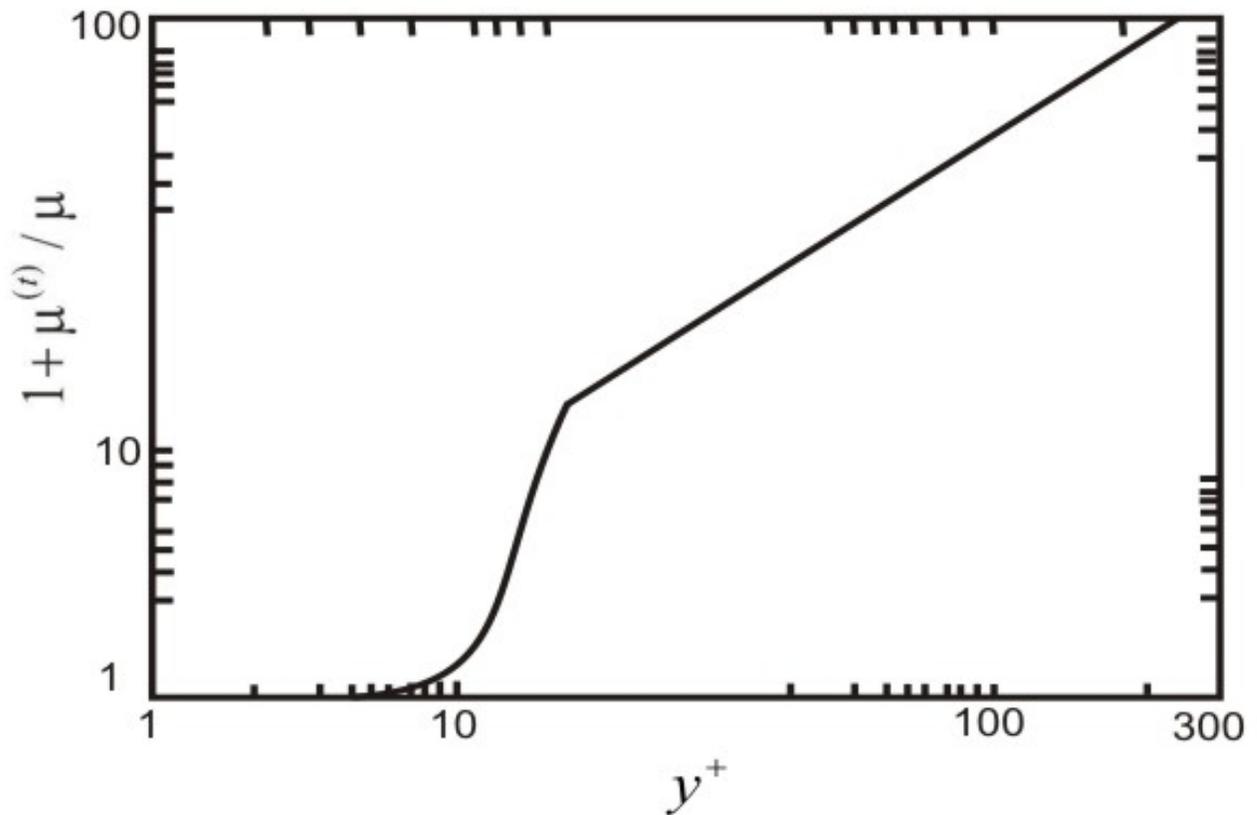
$$\bar{\tau}_{xy}^{(t)} = -\mu^{(t)} \frac{\partial \bar{u}_x}{\partial y}$$

Assuming that  $\bar{\tau}_{xy} \approx -\tau_0$  (i.e. not very far from wall)

Then,

$$\begin{aligned} \tau_0 &= (\mu + \mu^{(t)}) \frac{\partial \bar{u}_x}{\partial y} = \frac{\mu + \mu^{(t)}}{\mu} \tau_0 \frac{du^+}{dy^+} \\ \Rightarrow \left(1 + \frac{\mu^{(t)}}{\mu}\right) \frac{du^+}{dy^+} &= 1 \end{aligned} \quad (4.16)$$

$\Rightarrow \left(1 + \frac{\mu^{(t)}}{\mu}\right)$  is a universal function of  $y^+$ .



Measured velocity profiles do not give information for  $\mu^{(t)}$  near the wall since  $\mu^{(t)} \ll \mu$  in that region.

## Mass Transfer in Turbulent Flow

It is common to assume  $D^{(t)} = \nu^{(t)} = \frac{\mu^{(t)}}{\rho}$  (4.17)

and  $\bar{J}_{iy}^{(t)} = \text{turbulent mass flux} = \overline{-c'_i u'_y} = -D^{(t)} \frac{\partial \bar{c}_i}{\partial y}$

Near the wall ( $y^+ < 5$ )  $\Rightarrow \nu^{(t)} \ll \nu$

No info on  $\nu^{(t)}$  can be gathered very near wall.

But if  $Sc \gg 1$  (i.e.  $\nu \gg D_i$ ), it is necessary to know  $D^{(t)}$  closer to the wall. Thus, even if  $\nu^{(t)} = D^{(t)}$ ,

we can have at some distant  $D^{(t)} \gg D_i$ , even where  $\nu^{(t)} \ll \nu$ , if  $Sc \gg 1$ .

Thus much of our information about  $D^{(t)}$  (and possibly  $\nu^{(t)}$ ) near the wall comes from mass transfer experiments.

For example, in a pipe

$$St = \frac{D_i}{\langle \bar{u}_z \rangle \Delta c_i} \frac{\partial \bar{c}_i}{\partial r} \Big|_{r=R} = \frac{Sh}{Re Sc} \quad (4.18)$$

For  $Sc \gg 1$ , if  $D^{(t)} \sim y^n \Rightarrow St \sim Sc^{\frac{1-n}{n}}$ , where  $n = 2, 3, 4, \dots$

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