

## Module 2: "Capliarity"

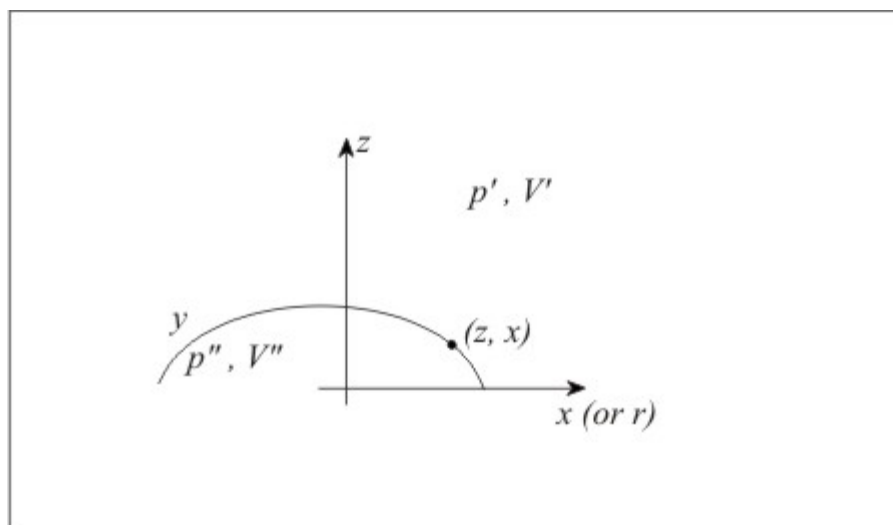
### Lecture 6: ""

The Lecture Contains:

- Axisymmetric Surfaces
- Young's Equation

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## Axisymmetric surfaces

Figure 2.6. Axisymmetric surface ( $z$  is the axis of symmetry)

An axisymmetric surface is a surface obtained by rotating a curve around an axis

For an axisymmetric surface (Fig. 2.6),

$$dA = 2\pi x dl = 2\pi x \sqrt{(dx)^2 + (dz)^2} \quad (2.22)$$

$$dV = (2\pi x dx)z \quad (2.23)$$

Consider a liquid drop on a surface (Figure 2.6).  $p', V'$ , and  $T$  being constant, the Helmholtz free energy ( $F$ ) is given as

$$F = U - TS + \sum_i \gamma_i A_i$$

Since,  $U = -pV + TS + G$

We get,

$$\therefore F = -pV + \sum_i n_i \mu_i + \sum_i \gamma_i A_i$$

$$F = \sum_i \gamma_i A_i - pV - p'V' + \text{const}$$

$$V' + V'' = V \Rightarrow V' = V - V''$$

$$F = \sum_i \gamma_i A_i - pV - p'(V - V'') + c$$

Denoting  $\Delta p = p' - p''$

$$F = \sum \gamma_i A_i - V'' \Delta p$$

$$= \gamma \int_0^R 2\pi x \left[ 1 + \left( \frac{dz}{dx} \right)^2 \right]^{1/2} dx + \int_0^R 2\pi x (\Delta p) z dx$$

$\Delta p$  also includes the body forces.

Here,  $z$  as a function of  $x$  is not known.

$F$  is a functional and variational calculus can be used to find its extrema.

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For a functional,

$$F = \int_0^{\kappa} f(x, z, z_x) dx \quad (2.24)$$

the condition for equilibrium is given by

$$\frac{\partial f}{\partial z} - \frac{d}{dx} \frac{\partial f}{\partial z_x} = 0 \quad (2.25)$$

In the present example,

$$F = \gamma \int_0^{\kappa} 2\pi x (1 + z_x^2)^{1/2} dx + \int_0^{\kappa} \Delta p 2\pi x z dx$$

$$\text{i.e. } f = \gamma 2\pi x (1 + z_x^2)^{1/2} + \Delta p 2\pi x z$$

So, for extrema

$$\Delta p = p' - p'' = \gamma \left[ \frac{z_x}{x(1 + z_x^2)^{1/2}} + \frac{z_{xx}}{x(1 + z_x^2)^{3/2}} \right] \quad (2.26)$$

The first term inside the braces is  $1/R_2$  while the other is  $1/R_1$ .

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This is the **Young-Laplace equation** for axisymmetric surfaces. It is a non-homogeneous second-order non-linear differential equation. This equation may have a unique solution, no solution or many solutions. A frequently used simplification is that for weakly curved surfaces,

$$\left(\frac{dz}{dx}\right)^2 \ll 1$$

### Some special cases

#### 1. Homogeneous Young-Laplace equation $\Delta p = 0$

One possible solution (trivial solution) is a flat surface situation, i.e.  $R_1$  &  $R_2$  both going to infinity. Another solution is a thin film formed between two loops kept at a distance away (Fig. 2.7).

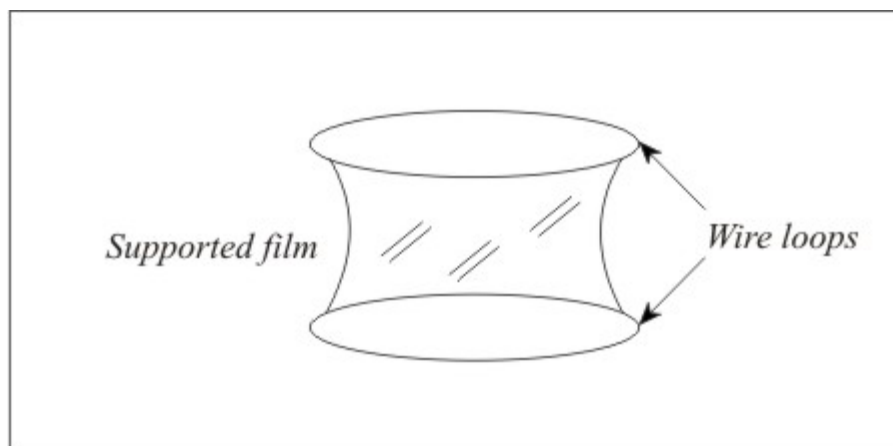


Figure 2.7. A thin film supported by two thin wire loops

#### 2. $\Delta p = \text{constant}$

Then  $R_1 = R_2 = R$ . This means that the surface is spherical.

## Young's Equation

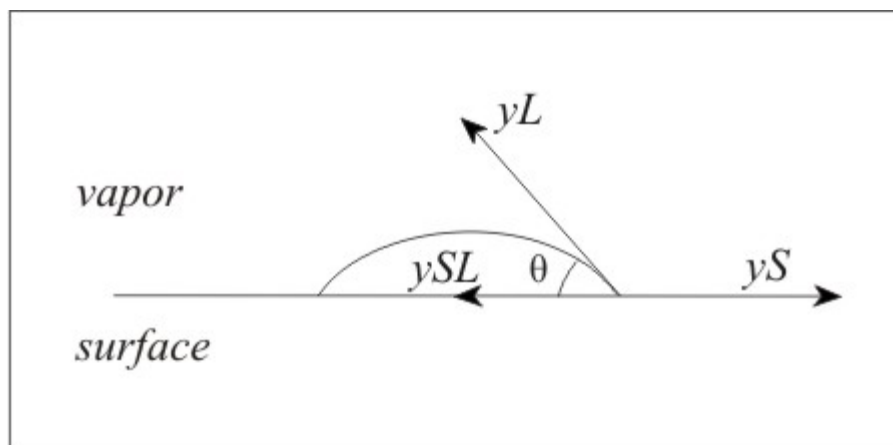


Figure 2.8. A liquid drop on a surface

$$F = \gamma_L \int_0^{\kappa} 2\pi x (1 + z_x^2)^{1/2} dx + \int_0^{\kappa} \Delta p 2\pi x z dx + \gamma_{SL} \int_0^{\kappa} 2\pi x dx + A\gamma_S - \gamma_S \int_0^{\kappa} 2\pi x dx$$

$$f = 2\pi x [\gamma_L (1 + Z_x^2)^{1/2} + \Delta p z + \gamma_{SL} - \gamma_S]$$

For equilibrium shape at the contact line

$$\left( f - z_x \frac{\partial f}{\partial z_x} \right)_{x=R} = 0$$

$$\gamma_L \frac{1}{(1 + z_x^2)^{1/2}} + \gamma_{SL} - \gamma_S = 0 \quad (2.27)$$

$$\gamma_L \cos \theta + \gamma_{SL} - \gamma_S = 0 \quad (2.28)$$

This is *Young's Equation*. It represents the balance of interfacial tension in the horizontal direction. Along the vertical, the interfacial tension is balanced by the elastic forces of the underlying surface.

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The contact angle is a measure of the wettability of the surface by that particular liquid. Higher the contact angle lesser is the wettability. The extreme situations correspond to  $\theta = 0^\circ$  when it is perfectly wettable and  $\theta = 180^\circ$  when it is non-wettable. A mercury drop is an example of a non-wettable liquid. Therefore, a mercury drop tries to form a perfect sphere but the sphere gets slightly flattened due to gravity.

The Young-Laplace and Young's equation considered together give us the shape of a surface. However, these equations are extremely difficult to solve analytically. Fortunately, various simplifications are possible in most cases which makes it much easier to solve the equations.

$$\gamma_{23} \cos \theta + \gamma_{12} = \gamma_{13} \tag{2.29}$$

$$\Delta P = \gamma \left[ \frac{z'}{(1 + z'^2)^{3/2}} + \frac{z'}{x(1 + z'^2)^{1/2}} \right] \tag{2.30}$$