

## Module 2: "Capliarity"

### Lecture 9: ""

The Lecture Contains:

- ☰ Hanging/Pendent Bubble
- ☰ Stability of water Cylinder (No Flow)

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## HANGING/PENDANT BUBBLE

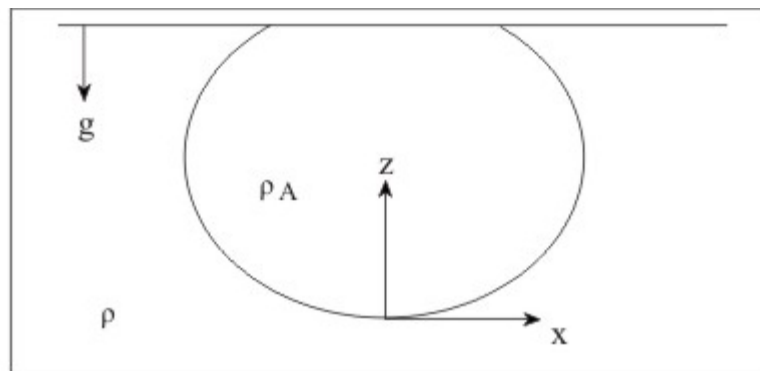


Figure:2.14 Pendant bubble

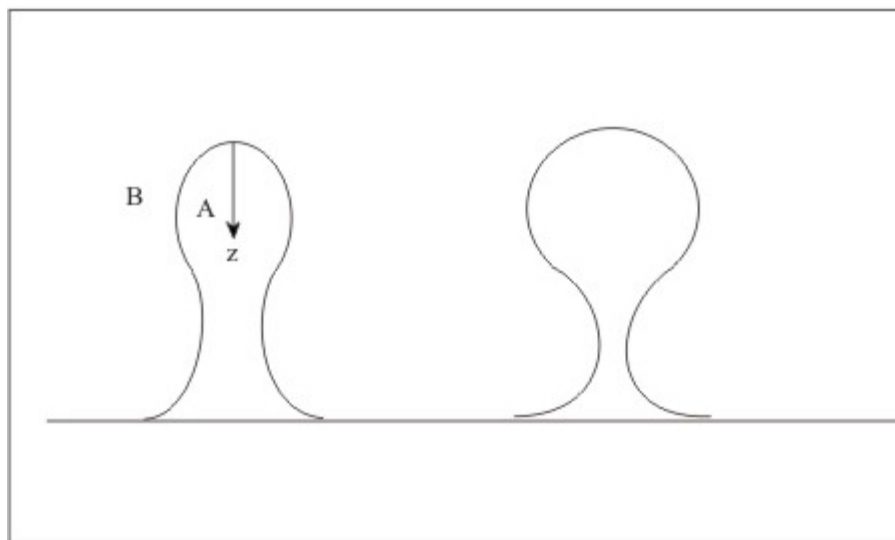


Figure:2.15. Necking

For a pendant bubble

$\rho_A < \rho_B$  (mirror image of the last case)

$$\Delta\rho = \rho_B - \rho_A > 0$$

Buoyancy would pull the drop upwards and a neck (fig:2.15) would form as the difference in density increases. If volume is more than a critical limit, the bubble would detach from the surface.

This phenomenon occurs only for contact angles  $\theta < 90^\circ$ ; hydrophobic surface.

Considering the case of a pendant drop, we take  $b$  as the length scale.

$$Z = \frac{z}{b}$$

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$$X = \frac{x}{b}$$

Using Young-Laplace equation,

$$2 + \left[ \frac{\Delta \rho g b^2}{\gamma} \right] Z = \left[ \frac{Z''}{(1 + Z'^2)^{3/2}} + \frac{Z'}{X(1 + Z'^2)^{1/2}} \right]$$

$$Z'' = \frac{d^2 Z}{dX^2}$$

$$Z' = \frac{dZ}{dX}$$

$$\left[ \frac{\Delta \rho g b^2}{\gamma} \right] \Rightarrow \text{Non dimensional number which is denoted by } (Ca)^{-1}$$

Where 'Ca' is called capillary number

If  $(Ca)^{-1} \ll 1$  then we can neglect the second term on the left hand side.

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## STABILITY OF A WATER CYLINDER (NO FLOW)

The surface solution will be in the form of

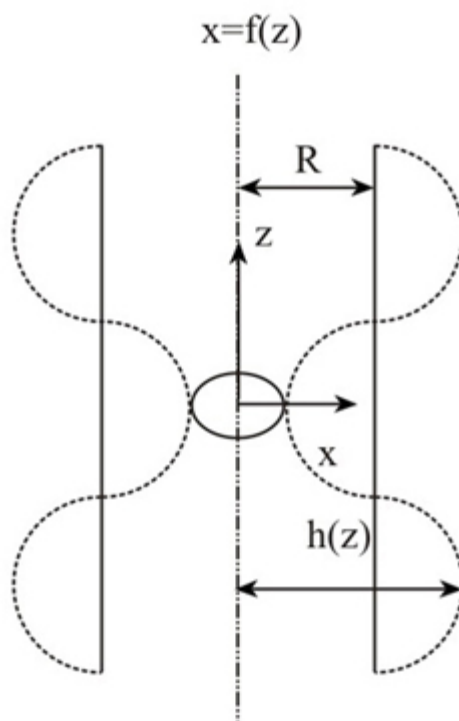


Fig 2.16

$$z = f(x)$$

For convenience we write the equation in the form  $x = f(z)$

$$\frac{d^2 z}{dx^2} = \frac{d}{dx} \left( \frac{df}{dz} \right)^{-1} = \frac{d}{dz} \left( \frac{df}{dz} \right)^{-1} \frac{dz}{dx} = - \left( \frac{df}{dz} \right)^{-2} \frac{d^2 f}{dz^2} \frac{dz}{dx} = - \left( \frac{df}{dz} \right)^{-3} \frac{d^2 f}{dz^2}$$

$$\frac{1}{R_1} = \frac{z''}{(1+z'^2)^{3/2}} = - \frac{x''}{(1+x'^2)^{3/2}}$$

$$\frac{1}{R_2} = \frac{z''}{x(1+z'^2)^{1/2}} = - \frac{x''}{x(1+x'^2)^{1/2}}$$

$$x = R + h(z)$$

$$\frac{1}{R_1} = - \frac{x''}{(1+x'^2)^{3/2}} = - \frac{h_{zz}}{(1+h_z^2)^{3/2}}$$

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if the slope is small

i.e.  $h_z \ll 1$

The scale on which a film is spread is very large in comparison to the thickness scale.

$$\frac{1}{R_1} = -h_{zz} \text{ (for small slope)}$$

$$\frac{1}{R_2} = \frac{1}{(R+h)(1+x'^2)^{1/2}} \rightarrow 0 \text{ if } R \rightarrow \infty$$

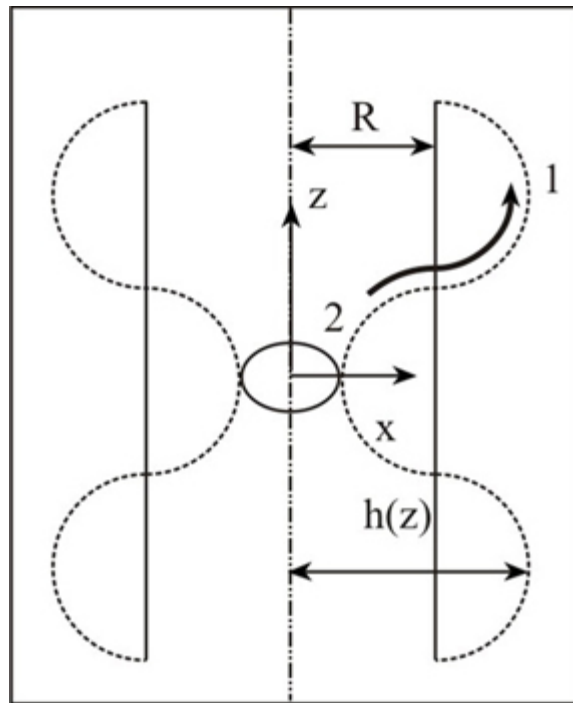


Fig 2.17:Disturbance of a water cylinder