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Hartree Fock formalism

PCD_STiAP_P04

1. For a two electron system, spin exchange operator is given by $\hat{V}_{ex} = -\frac{1}{2}K(1 + 4\hat{s}_1 \cdot \hat{s}_2)$

Where K is the exchange integral and \hat{s}_1 and \hat{s}_2 are spin angular momentum operators of two electrons. This was first introduced by Dirac and has wide applications in Condensed Matter Physics.

a) Show that $\hat{s}_1 \cdot \hat{s}_2 = \frac{1}{2}[\hat{S}^2 - \hat{s}_1^2 - \hat{s}_2^2]$

Where \hat{S} is total spin operator.

- b) Show that eigenvalues of $\hat{s}_1 \cdot \hat{s}_2$ for triplet ($S = 1$) and singlet ($S = 0$) states are $\frac{1}{4}$ and $-\frac{3}{4}$ respectively.
- c) Find the eigenvalues of V_{ex} for triplet and singlet states.
- d) Show that $V_{singlet} - V_{triplet} = 2K$.
- e) Which is the ground state of the system?

2. Prove that the direct integral J and exchange integral K for a N electron atomic system are real and positive.
3. For the ground state of an atom or ion having N electrons,

$$H_1 = \sum_{i=1}^N f(q_i) = \sum_{i=1}^N \left(-\frac{\nabla_i^2}{2} - \frac{Z}{r_i} \right)$$

$$H_2 = \frac{1}{2} \sum_{i,j=1, j \neq i}^N V(q_i, q_j) = \sum_{i>j} \frac{1}{r_{ij}}$$

Where H_1 is the sum of the N identical one-body hydrogenic Hamiltonian and H_2 is the sum of $\frac{N(N-1)}{2}$ identical terms which represent the two-body interactions between each pair of electrons and ϕ_a and ϕ_b are Slater determinants,

$$\phi_a(q_1, q_2, \dots, q_N) = \frac{1}{N!} \sum_P (-1)^P P u_{\alpha_1}(q_1) u_{\alpha_2}(q_2) \dots u_{\alpha_N}(q_N)$$

$$\phi_b(q_1, q_2, \dots, q_N) = \frac{1}{N!} \sum_P (-1)^P P u_{\beta_1}(q_1) u_{\beta_2}(q_2) \dots u_{\beta_N}(q_N)$$

where P is the permutation operator.

Prove that:

a) For $\phi_a = \phi_b$,

$$\begin{aligned}\langle \phi_a | H_1 | \phi_a \rangle &= \sum_{i=1}^N \langle u_{\alpha_i}(q_i) | f | u_{\alpha_i}(q_i) \rangle \\ \langle \phi_a | H_2 | \phi_a \rangle &= \frac{1}{2} \sum_{i,j=1}^N [\langle u_{\alpha_i}(q_i) u_{\alpha_j}(q_j) | V | u_{\alpha_i}(q_i) u_{\alpha_j}(q_j) \rangle \\ &\quad - \langle u_{\alpha_j}(q_i) u_{\alpha_i}(q_j) | V | u_{\alpha_i}(q_i) u_{\alpha_j}(q_j) \rangle].\end{aligned}$$

b) For $\phi_a \neq \phi_b$

($\beta_k \neq \alpha_k$ and $\beta_i = \alpha_i$ for all $i \neq k$)

$$\begin{aligned}\langle \phi_b | H_1 | \phi_a \rangle &= \langle u_{\beta_k}(q_k) | f | u_{\alpha_k}(q_k) \rangle \\ \langle \phi_b | H_2 | \phi_a \rangle &= \sum_{i=1}^N \langle u_{\alpha_i}(q_i) u_{\beta_k}(q_k) | V | u_{\alpha_i}(q_i) u_{\alpha_k}(q_k) \rangle \\ &\quad - \langle u_{\beta_k}(q_i) u_{\alpha_i}(q_k) | V | u_{\alpha_i}(q_i) u_{\alpha_k}(q_k) \rangle\end{aligned}$$

c) For $\phi_a \neq \phi_b$

($\beta_k \neq \alpha_k; \beta_k \neq \alpha_l; \beta_l \neq \alpha_l; \beta_l \neq \alpha_k$ and $\beta_i = \alpha_i$ for all $i \neq k, l$)

$$\begin{aligned}\langle \phi_b | H_1 | \phi_a \rangle &= 0 \\ \langle \phi_b | H_2 | \phi_a \rangle &= \sum_{i=1}^N \langle u_{\beta_k}(q_k) u_{\beta_l}(q_l) | V | u_{\alpha_k}(q_k) u_{\alpha_l}(q_l) \rangle \\ &\quad - \langle u_{\beta_l}(q_k) u_{\beta_k}(q_l) | V | u_{\alpha_k}(q_k) u_{\alpha_l}(q_l) \rangle.\end{aligned}$$

d) For $\phi_a \neq \phi_b$

(Differing in more than 2 sets of q labels.)

$$\begin{aligned}\langle \phi_b | H_1 | \phi_a \rangle &= 0 \\ \langle \phi_b | H_2 | \phi_a \rangle &= 0\end{aligned}$$

4. Prove that the matrix λ , made up of Lagrange variational multipliers which constraint the variation in the elements of Hartree-Fock Slater determinants, is self adjoint.

5.

a) Obtain the condition that $\langle H \rangle$, where H is the N-electron Hamiltonian, is an extremum subject to the constraint that the one-electron spin-orbitals in the variational N-electron antisymmetrized wave function are normalized and orthogonal.

b) Express the above condition in a form in which the matrix of the Lagrange variational multipliers is diagonal and demonstrate that this condition is expressed by:

$$\begin{aligned}
f(\vec{r}_1)u_k(\vec{r}_1) + \left[\sum_j \int d\tau_2 \frac{u_j^*(\vec{r}_2)u_j(\vec{r}_2)}{r_{12}} \right] u_k(\vec{r}_1) \\
- \sum_j \delta(m_{s_k}, m_{s_j}) \left[\int d\tau_2 \frac{u_j^*(\vec{r}_2)u_k(\vec{r}_2)}{r_{12}} \right] u_j(\vec{r}_1) \\
= -\lambda_{kk}u_k(\vec{r}_1) \\
= \varepsilon_k u_k(\vec{r}_1)
\end{aligned}$$

Where $\varepsilon_k = -\lambda_{kk}$

c) Comment on the factor $\delta(m_{s_k}, m_{s_j})$ in the above equation.

d) Find the function $U_k(\vec{r}_1, \vec{r}_2)$ such that;

$$\int d\tau_2 U_k(\vec{r}_1, \vec{r}_2) u_k(\vec{r}_2) = - \sum_j \delta(m_{s_k}, m_{s_j}) \left[\int d\tau_2 \frac{u_j^*(\vec{r}_2)u_k(\vec{r}_2)}{r_{12}} \right] u_j(\vec{r}_1)$$

e) Show that the one-electron Hartree-Fock equation is written in a form inclusive of the spin variables such that the spin-orbital $U_k(q) = \langle q | k \rangle$ has k which represents a set of 4 quantum numbers including spin is given by:

$$\left[\frac{-\nabla_i^2}{2} + V_{HF}(q_i) \right] u_k(q_i) = \varepsilon_k u_k(q_i)$$

Where

$$V_{HF}(q_i) = \frac{-Z}{r_i} + V^c(q_i) + V^{ex}(q_i)$$

$$V^c(q_i) = \sum_{j=1}^N \int dq_2 \frac{u_j^*(q_2)u_j(q_2)}{r_{i2}}$$

$$V^{ex}u_k(q_i) = - \sum_{j=1}^N \int dq_2 \frac{u_j^*(q_2)u_k(q_2)}{r_{i2}} u_j(q_i)$$

f) Is the Hartree-Fock equation given above an eigenvalues equation?

g) Can you define the Hartree-Fock potential V_{HF} as a function of a single coordinate (that is, one set of 3 space coordinates and 1 spin coordinate) alone?

6. Determine if the potential U in problem 5(d) is Hermitian.

7.

a) Prove that;

$$-\lambda_{kk} = n_k \langle k | f | k \rangle + n_k \sum_j n_j [\langle kj | V | kj \rangle - \langle jk | V | kj \rangle]$$

b) If $-\lambda_{kk}$ is written as ε_k , prove that:

$$E(\phi^{(N)}) - E(\phi_{n_k=0}^{(N-1)}) = \varepsilon_k$$

[Ref : T.H. Koopmans Physics **1 104** (1933)]

d) Write the Slater determinant for ground state $1s^2 2s^2 {}^1S$ of beryllium.

e) The Hartree-Fock potential is given by:

$$V_{HF} = -\frac{4}{r_i} + V_{1s\uparrow}^d + V_{1s\downarrow}^d + V_{2s\uparrow}^d + V_{2s\downarrow}^d - (V_{1s\uparrow}^{ex} + V_{1s\downarrow}^{ex} + V_{2s\uparrow}^{ex} + V_{2s\downarrow}^{ex})$$

Where V^d and V^{ex} are direct and exchange integrals respectively.

Obtain the two coupled integro-differential equation:

$$\left\{ -\frac{1}{2} \nabla_r^2 - \frac{4}{r} + V_{1s}^d(r) + 2V_{2s}^d(r) - V_{2s}^{ex}(r) \right\} u_{1s}(r) = E_{1s} u_{1s}(r)$$

$$\left\{ -\frac{1}{2} \nabla_r^2 - \frac{4}{r} + V_{2s}^d(r) + 2V_{1s}^d(r) - V_{1s}^{ex}(r) \right\} u_{2s}(r) = E_{2s} u_{2s}(r)$$

8. Prove that $[H, L] = 0$, where H is the Hamiltonian

$$H = \sum_{i=1}^N \left(-\frac{\nabla_i^2}{2} - \frac{Z}{r_i} \right) + \sum_{i < j=1}^N \frac{1}{r_{ij}}$$

And $L = \sum_i L_i$ is the total orbital angular momentum of the electrons.

Useful references:

1. Bethe, H.A and Jackiw R; Intermediate Quantum Mechanics; West-view press (1986).
2. Brandsen aB.H and joachain C.J; Physics of atoms and molecules; Longman Group Limited (1983).
3. Landau L.D. and Lifshitz E.M; Quantum Mechanics Non-Relativistic Theory; Pergamon Oxford (1977)