

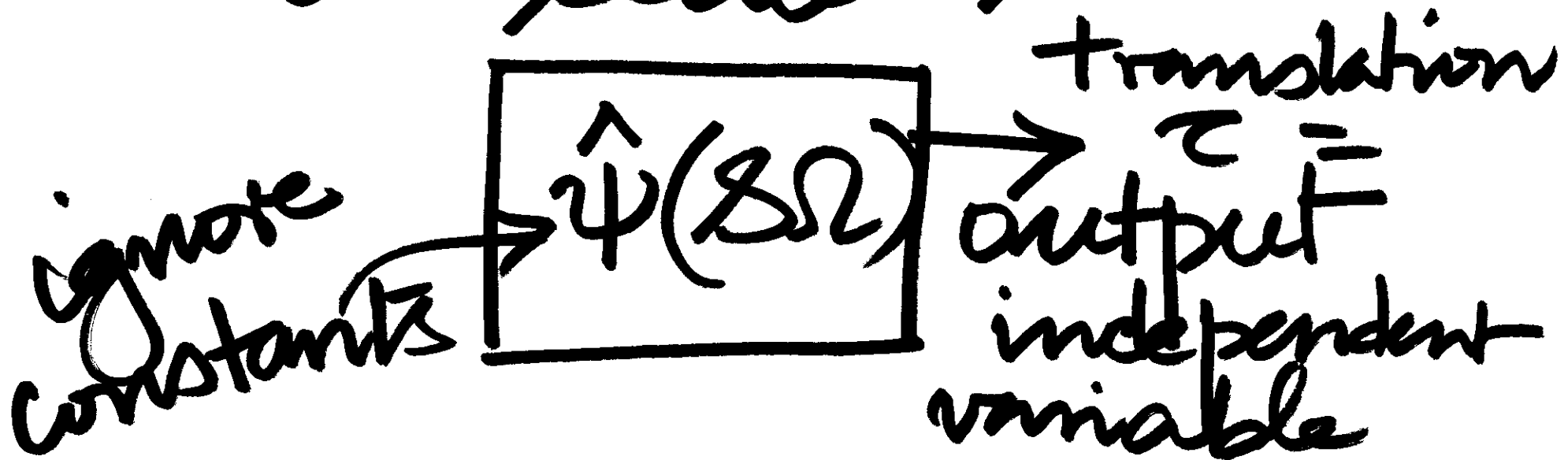
# LECTURE 24

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Lec. No. 24  
Date: 25/2/11

LOGARITHMIC SCALE  
DISCRETIZATION  
DYADIC DISCRETIZATION

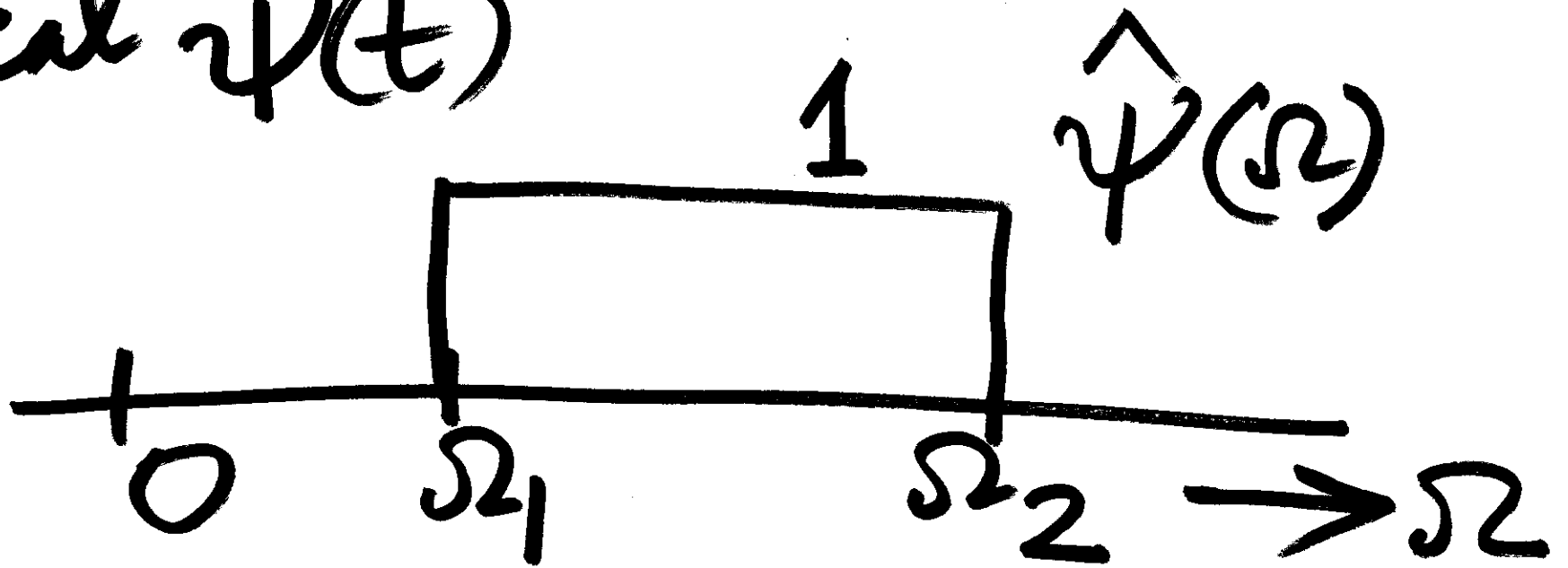
# Continuous Wavelet Transform

at scale  $\mathcal{S}$



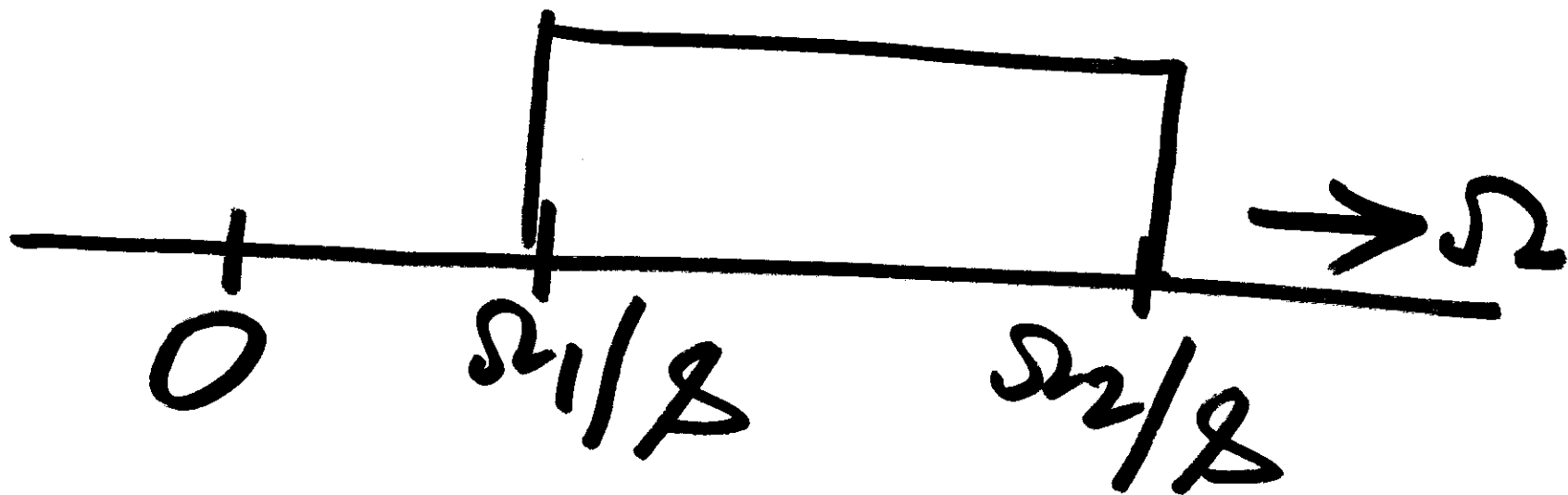
'Ideal' admissible

'wavelets':  
real  $\psi(t)$



For any  $\delta > 0$

$$\hat{\psi}(\omega)$$

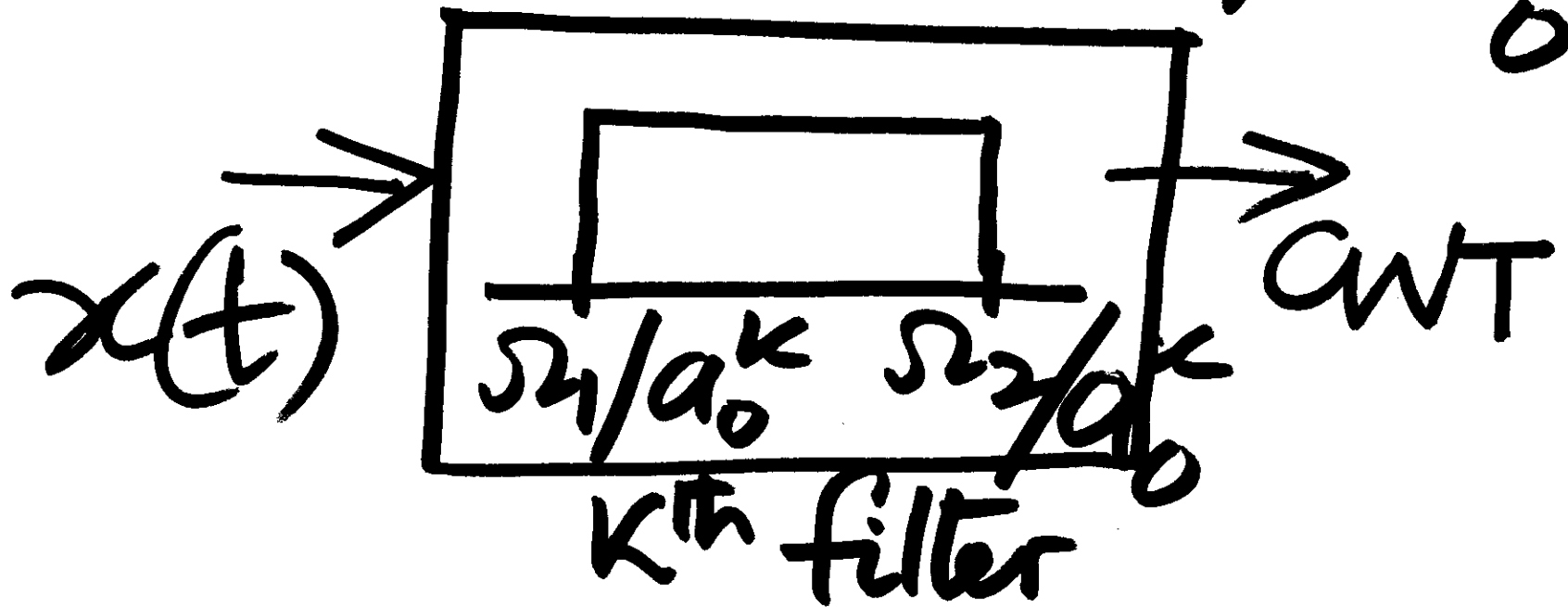


logarithmic  
discretization

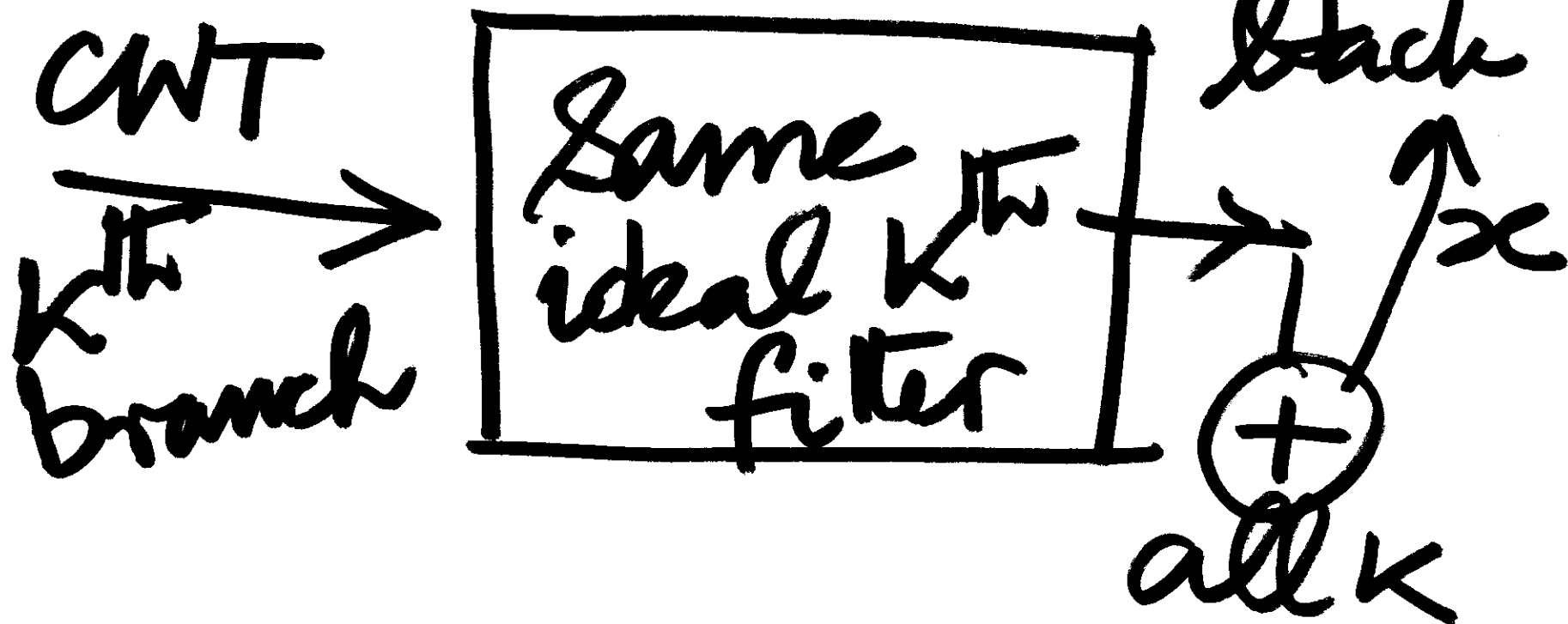
$$\Rightarrow \beta = a_0^k, \quad \underline{a_0 > 1}$$

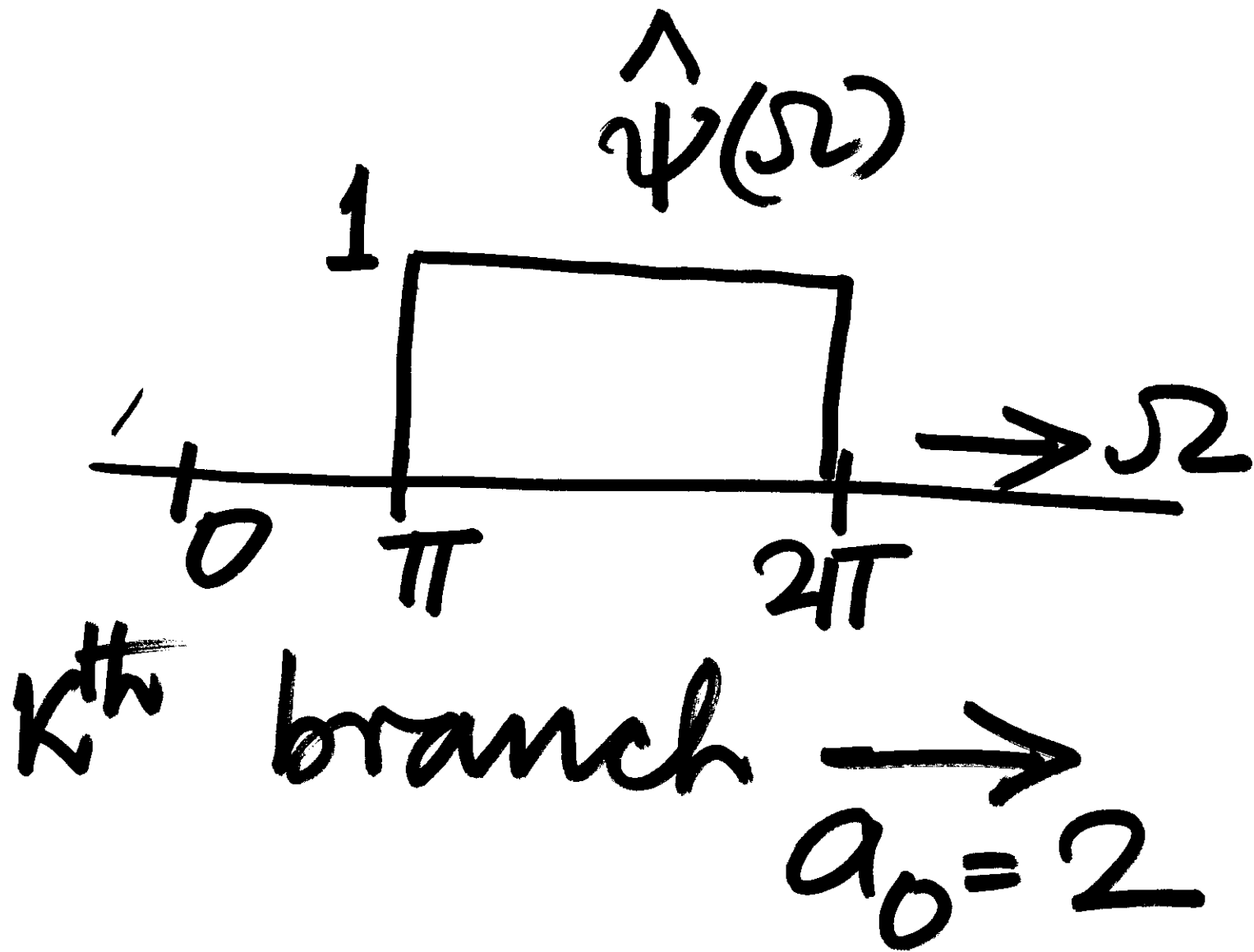
k: all integer values

For each  $k$  we have a filter  $\beta = a_0^k$

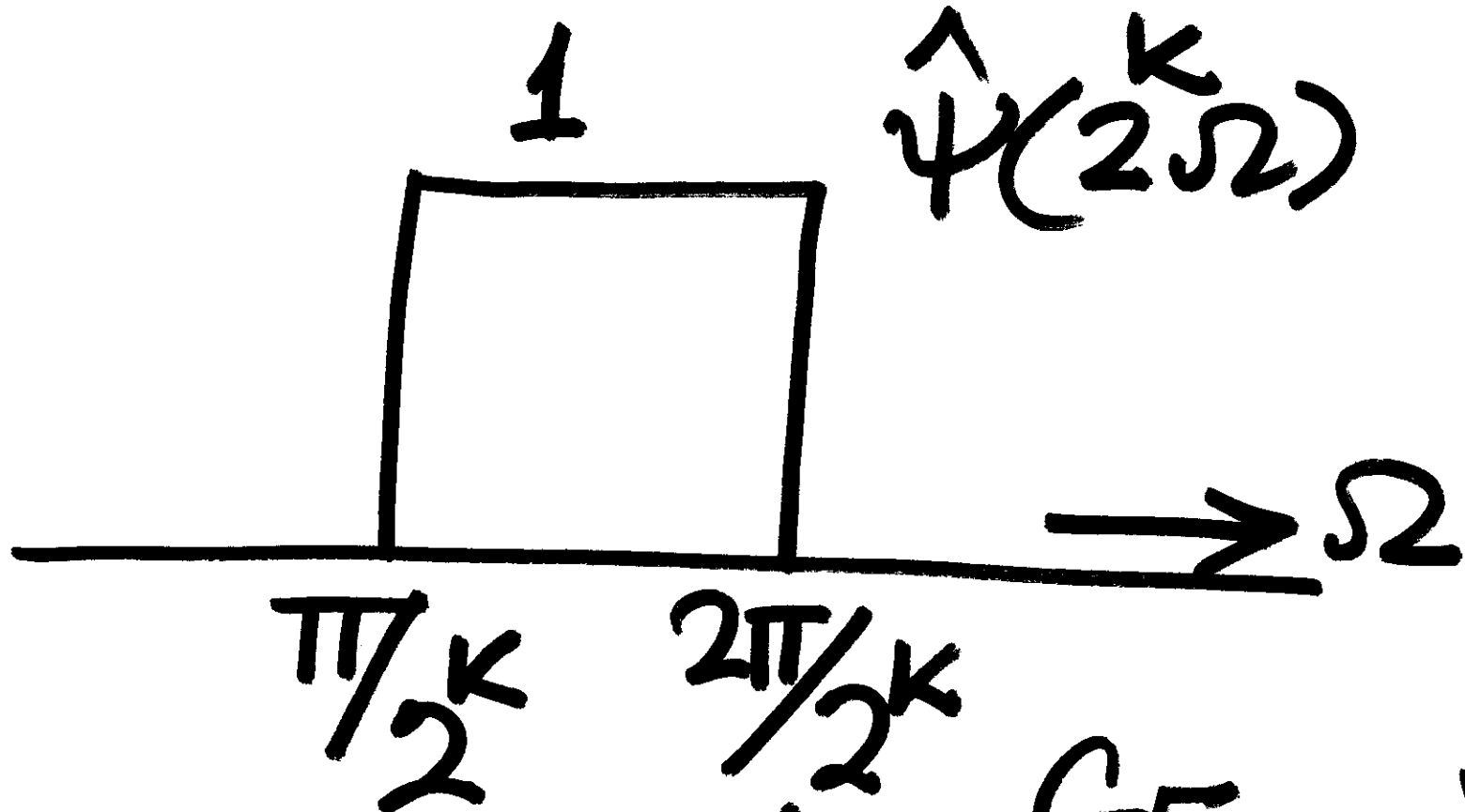


$k^{\text{th}}$  branch of reconstruction





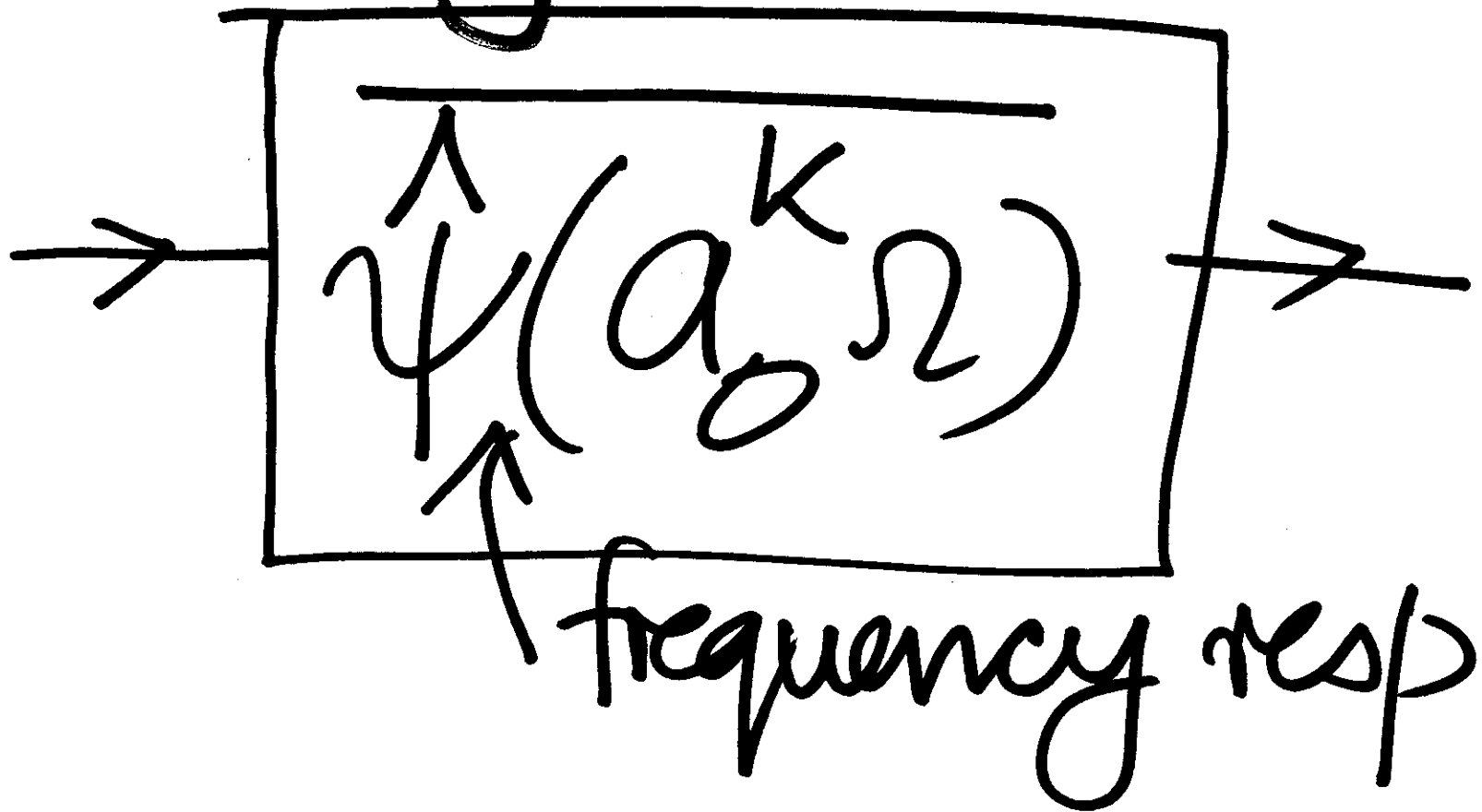




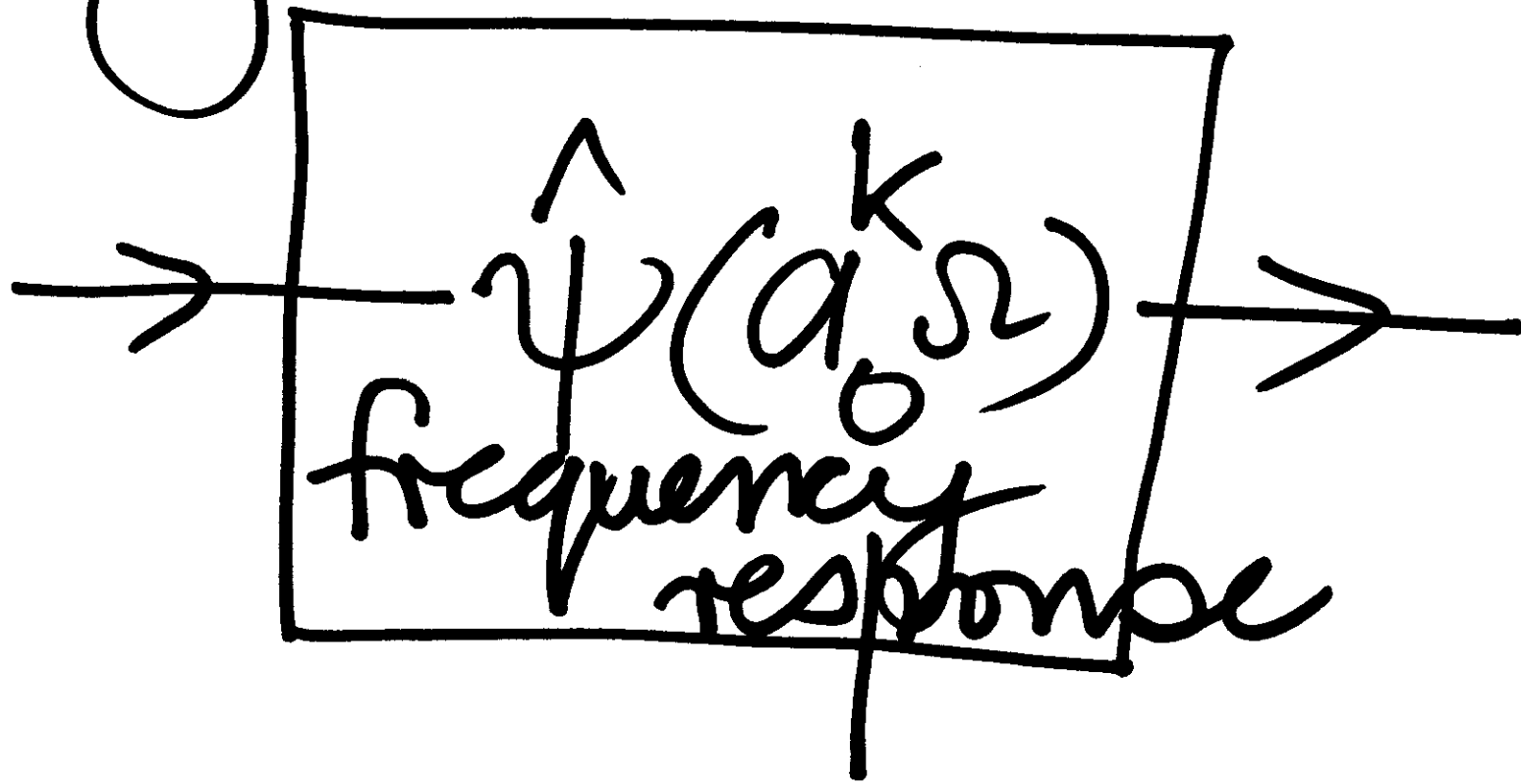
Nonoverlapping for integer  $k$

Discretization of scale  
parameter:  
equivalent to  
Constructing filter bank

# "Analysis" filters



# 'Synthesis' filters



On the  $k^{\text{th}}$  branch

$$\begin{array}{ccc} x(t) & \xrightarrow{\text{Fourier}} & \hat{x}(\Omega) \\ \text{input} & & \\ \downarrow & & \\ \hat{x}(\Omega) & \hat{\psi}(a_0^k \Omega) & \end{array}$$

Transform ANALYSIS

$k^{\text{th}}$  Synthesis branch:

$$\begin{aligned} & \overline{\hat{x}(\omega) \hat{\psi}(a_0^k \omega) \hat{\psi}(a_0^k \omega)} \\ &= \hat{x}(\omega) |\hat{\psi}(a_0^k \omega)|^2 \end{aligned}$$

Overall output

$$= \hat{x}(\omega) \left\{ \sum_{k=-\infty}^{+\infty} |\hat{\psi}(a_0^k \omega)|^2 \right\}$$

Ideally, this  $\uparrow$   $= \frac{1}{T\omega}$

To relax this requirement for  
'designability':

$$C_1 \leq \sum_{k=-\infty}^{+\infty} |\hat{\psi}(a_0, \omega^k)|^2 \leq C_2$$

$$0 < C_1 \leq C_2 < +\infty$$



In that case,

define another

$$\tilde{\psi}(t) \xrightarrow{\text{Fourier Transform}} \tilde{\psi}(\omega)$$

$$\hat{\psi}^2(\Omega) = \hat{\psi}(\Omega)$$

$$\frac{+\infty}{\sum_{k=-\infty}^{+\infty} |\hat{\psi}(a_0^k \Omega)|^2}$$

Valid because of  $\Omega > 0$ .

We first show, ( $\psi$  real)  
 $\psi$  is automatically  
admissible.

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$$\int_0^{\infty} |\hat{\psi}(\alpha)|^2 \frac{d\alpha}{\alpha} \quad ?$$

$$\int_0^{\infty} |\hat{\psi}(\alpha)|^2 \frac{d\alpha}{\alpha}$$

$$= \sum_{k=-\infty}^{\infty} a_k$$

$$a_k \xrightarrow{a_{kH}} a_0^k$$

$$\int_0^{\infty} |\hat{\psi}(\alpha)|^2 \frac{d\alpha}{\alpha} \text{ put } \alpha = a_0^k \beta$$

$$\int_{a_0^k}^{a_0^{k+1}} |\hat{\psi}(\alpha)|^2 \frac{d\alpha}{\alpha} \quad d\alpha = a_0^k d\beta \\
 \alpha = a_0 \beta \\
 = \int_1^{a_0} \hat{\psi}\left(\frac{a_0^k}{\beta}\right) \frac{d\beta}{\beta}$$

$$\int_0^{\infty} |\hat{\psi}(\alpha)|^2 \frac{d\alpha}{\alpha}$$

$$= \int_1^{\infty} \left\{ \sum_{k=-\infty}^{+\infty} |\hat{\psi}(a_0^k \beta)|^2 \right\} \frac{d\beta}{\beta}$$

$$\sum_{k=-\infty}^{+\infty} |\hat{\psi}(a_0^k \beta)|^2$$

is upper bounded  
by  $C_2$

'Admissibility  
integral'  $\int_0^{\infty} |\psi(x)|^2 dx$   
is upper bounded:  
by



$$\int_1^{a_0} C_2 \frac{d\beta}{\beta} = C_2 \ln \beta \Big|_1^{a_0} = C_2 \ln a_0$$

$$\sum_{k=-\infty}^{+\infty} |\hat{\psi}(a_0^k \Omega)|^2$$

Sum of dilated  
spectra (SDS)

SDS  $(\psi, a_0)$   $(\Omega)$

Secondary primary  
arguments

~~SDS(a\_0)~~

SDS( $\psi, a_0$ ) ( $\Omega$ )

$C_1 \searrow$

$\searrow C_2$

$C_2$  guarantees

ADMISSIBILITY.

$$\hat{\tilde{\psi}}(\Omega) = \frac{\hat{\psi}(\Omega)}{\text{SDS}(\psi, a_0)(\Omega)}$$

We show that  $\tilde{\psi}$  is  
ADMISSIBLE

For this purpose we  
need to consider

$$SDS(\tilde{\psi}, a_0)(\Omega)$$

$$|\hat{\psi}(a_0^k \Omega)|^2 = |\hat{\psi}(a_0^k \Omega)|^2$$

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$$\left\{ \sum_{l=-\infty}^{+\infty} |\hat{\psi}(a_0^l a_0^k \Omega)|^2 \right\}$$

$$|\hat{\psi}(a_0^k, \Omega)|^2 = |\hat{\psi}(a_0^k, \Omega)|^2$$

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$$\{SDS(\psi, a_0)(\Omega)\}^2$$



$$\text{SDS}(\tilde{\psi}, a_0)(\Omega)$$

$$= \sum_{k=-\infty}^{+\infty} |\hat{\tilde{\psi}}(a_0^k \Omega)|^2$$

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$$\{\text{SDS}(\psi, a_0)(\Omega)\}^2$$

$$SDS(\tilde{\psi}, a_0)(\Omega)$$

$$= \frac{SDS(\psi, a_0)(\Omega)}{\{SDS(\psi, a_0)(\Omega)\}^2}$$

$$\{SDS(\psi, a_0)(\Omega)\}^2$$

Cancellation from  
numerator and  
denominator is  
valid because of  
bounds  $c_1, c_2$

$SDS(\tilde{\psi}, a_0)(\Omega)$ 

$$= \frac{1}{SDS(\psi, a_0)(\Omega)}$$

$$0 < c_1 \leq \text{SDS}(\psi, a_0)(\Omega) \leq c_2 < \infty$$



$$\infty > \frac{1}{c_1} \geq \text{SDS}(\psi, a_0)(\Omega) \geq \frac{1}{c_2} > 0$$

$\text{SDS}(\psi, a_0)(\Omega)$

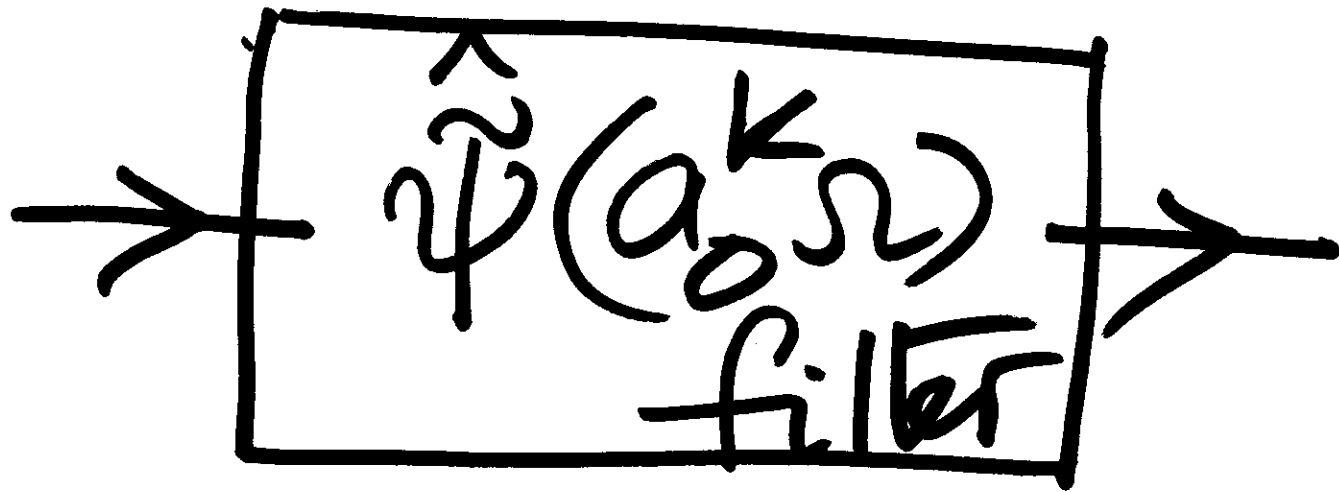
$\psi^2$  is also an

ADMISSIBLE

WAVELET

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Use  $\hat{\psi}$  on synthesis  
side:  $k^{\text{th}}$  branch:



With this, analysis  
and synthesis  
together:

$$\hat{x}(\omega) \underset{k}{=} \hat{\psi}(a_0^k \omega) \cdot \hat{\tilde{\psi}}(a_0^k \omega)$$



$$\hat{x}(\Omega) \cdot \sum_{k=-\infty}^{+\infty} \dots$$

$$\dots \frac{\hat{\psi}(a_0^k \Omega)}{\text{SDS}(\psi, a_0)(\Omega)}$$

$$= \hat{x}(\Omega) \cdot \frac{\text{SDS}(\psi, a_0)(\Omega)}{\text{SDS}(\psi, a_0)(\Omega)}$$

Because of  $c_1, c_2$   
we can cancel  $\psi$

$$= \hat{x}(\Omega) \times 1$$

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Specific case  
of  $a_0 = 2$   
Dyadic'

$SDS(4, 2)(\Omega)$

must be upper  
and lower  
bounded.

Exercise: For the  
Haar wavelet  $\psi(t)$   
show that  
 $0 < C_1 \leq SDS(\psi, 2)(\Omega) \leq C_2 < \infty$

For the "Derivative  
of Gaussian"  
wavelets, a very  
wide range of  $a_0$  is possible.