

LECTURE 23

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Lec-23
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ADMISSIBILITY IN DETAIL DISCRETIZATION OF SCALE

Important steps in
reconstructing $x(t)$
from $CWT(x, \psi)(\tau, \lambda)$

Triple integral involved

$$\frac{1}{2\pi} \int \hat{x}(\Omega) \cdot \delta \cdot \overline{\hat{\psi}(\Omega)} e^{j\Omega\tau} d\Omega$$

$\underbrace{\hspace{15em}}$

$$\text{CWT}(x, \psi)(\tau, \delta)$$

Reconstruct:

$$\int_0^{\infty} \int_{-\infty}^{+\infty} \text{CWT}(x, \nu)(\tau, \delta) \frac{1}{\delta^{1/2}} \psi\left(\frac{t-\tau}{\delta}\right) f(\delta) d\tau d\delta$$

Writing all together

Overall integral:

$$\int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \hat{x}(\omega) \cdot \frac{1}{s} \cdot \hat{\psi}(s\omega) e^{j\omega\tau}$$

$$\cdot \frac{1}{s^{1/2}} \cdot \psi\left(\frac{t-\tau}{s}\right) f(s)$$

take last $\rightarrow \frac{d\omega}{\text{first}} \frac{d\tau}{\text{second}} \frac{ds}{\text{second}}$

After taking care of
 $\int_{-\infty}^{+\infty} dz$

what was left, by
choosing $f(z) = \frac{1}{z^2}$

$$\frac{1}{2\pi} \int_{-\infty}^{+\infty} \hat{x}(\Omega) \left\{ \int_0^{\infty} |\hat{\psi}(\Omega)|^2 \frac{ds}{s} \right\} e^{j\Omega t} d\Omega$$

If we could make

$$\int_0^{\infty} |\hat{\psi}(s\Omega)|^2 \frac{ds}{s}$$

independent
of Ω

(constant): C_{ψ}

What is left is:

$$\frac{1}{2\pi} \cdot C_y \int_{-\infty}^{+\infty} \hat{x}(\omega) e^{j\omega t} d\omega$$

$$= C_y \cdot x(t)$$

Consider: $\int_0^{\infty} |\hat{\psi}(z)|^2 \frac{dz}{z}$

Put $z = \alpha$

$z: 0 \rightarrow +\infty$

$\Omega = 0$ is a problem
Except for $\Omega = 0$,

$$\Omega > 0: \alpha = 2\Omega$$

$$\alpha: 0 \longrightarrow +\infty$$

when $\beta: 0 \longrightarrow +\infty$

$$d\alpha = \Omega d\beta$$

$$\alpha = \Omega \beta$$

$$\Omega \neq 0$$

$$\frac{d\alpha}{\alpha} = \frac{d\beta}{\beta}$$

$$\int_0^{\infty} |\hat{\psi}(\omega)|^2 \frac{d\omega}{\omega}$$

$$= \int_0^{\infty} |\hat{\psi}(\alpha)|^2 \frac{d\alpha}{\alpha}$$

If $\psi(t)$ is a
complex wavelet,
we need separately
to take care of the
positive and negative
spectrum

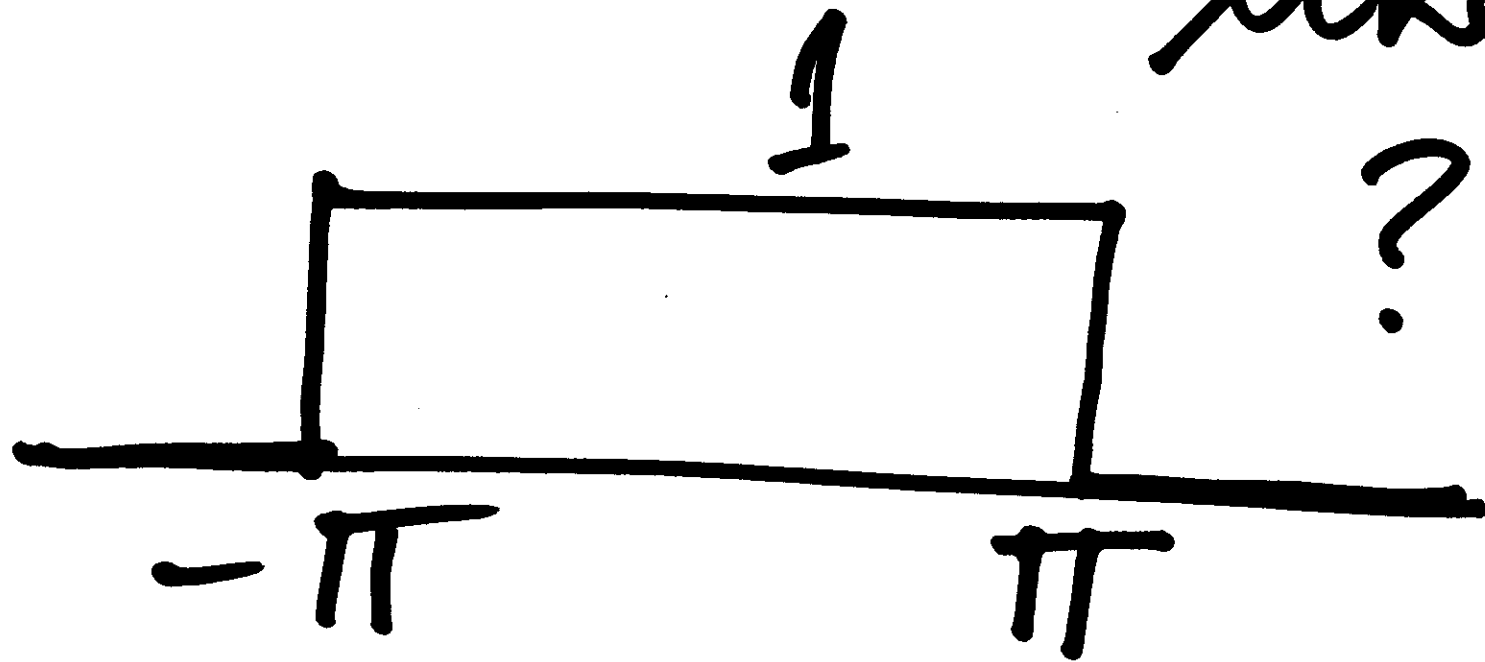
$$\int_0^{\infty} |\hat{\psi}(\alpha)|^2 \frac{d\alpha}{\alpha} \quad \text{ADMISSIBILITY}$$

$$= \int_0^{\infty} |\hat{\psi}(-\beta)|^2 \frac{d\beta}{\beta} < \infty.$$

'Admissibility'
is a condition
required for
RECONSTRUCTION .

$\Omega > 0$ Condition
(For real ψ this is
enough)
 $\int_0^\infty |\hat{\psi}(\alpha)|^2 \frac{d\alpha}{\alpha}$ finite

Can $|\hat{\psi}(x)|^2$ look like:



$$\Omega < 0: \quad \alpha = 2\Omega$$

$$\beta: 0 \longrightarrow +\infty$$

$$\alpha: 0 \longrightarrow -\infty$$

$$\int_0^{-\infty} |\hat{\psi}(\alpha)|^2 \frac{d\alpha}{\alpha}$$

Substitute $\beta = -\alpha$
 $\alpha: 0 \rightarrow \text{~~+\infty~~}; \beta: 0 \rightarrow +\infty$
 $\text{~~-\infty~~}$

$$\int_0^{\infty} |\hat{\chi}(\beta)|^2 \frac{d\beta}{\beta} \cdot \text{for } \underline{\Omega < 0}$$

$$\frac{d\beta}{\beta} = - \frac{d\alpha}{\alpha} \Rightarrow \frac{d\beta}{\beta} = \frac{d\alpha}{\alpha}$$

Making this integral
(only)
independent of Ω ,
means:

If $\psi(t)$ is real

$$\hat{\psi}(-\beta) = \overline{\hat{\psi}(\beta)}$$

$$|\hat{\psi}(-\beta)|^2 = |\hat{\psi}(\beta)|^2$$

$$\rho \int_0^{\infty} |\hat{\psi}(\alpha)|^2 \frac{d\alpha}{\alpha}$$

$$= \pi \int_0^{\infty} 1 \cdot \frac{d\alpha}{\alpha}$$

Divergent!

$$\int \frac{dx}{x} \ln x \Big|_0^\pi$$

divergent

Here the trouble
comes as $\alpha \rightarrow 0$
The spectrum
should have vanished
as $\alpha \rightarrow 0$

What about



$$|\hat{\psi}(\infty)|^2 \xrightarrow{\Omega} c_0 \text{ as } \alpha \rightarrow +\infty$$

$$\int_0^L |\hat{\psi}(x)|^2 \frac{dx}{L}$$

$L \rightarrow$ large enough

$$\approx \int_0^L c_0 \frac{dx}{L}$$

$\ln \alpha$ $\left(\begin{array}{c} f\infty \\ L \end{array} \right)$

divergent again!

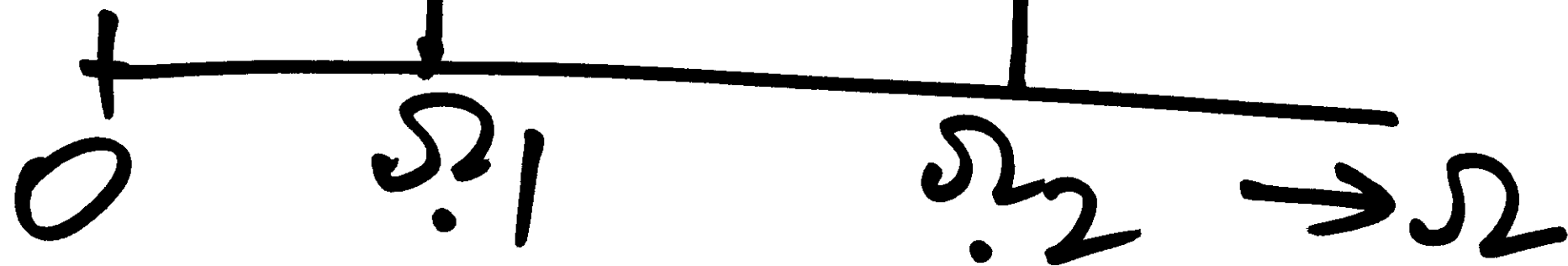
The spectrum $|\hat{\psi}(\alpha)|^2$
should have
vanished
as $\alpha \rightarrow +\infty$

Can we allow:

BANDPASS

1

$|\hat{\psi}(\omega)|^2$



$$\begin{aligned}
 & \int_0^{\infty} |\hat{\psi}(\alpha)|^2 \frac{d\alpha}{\alpha} \\
 &= \int_{\Omega_1}^{\Omega_2} 1 \cdot \frac{d\alpha}{\alpha}
 \end{aligned}$$

$$= \cancel{\ln \Omega_2}$$

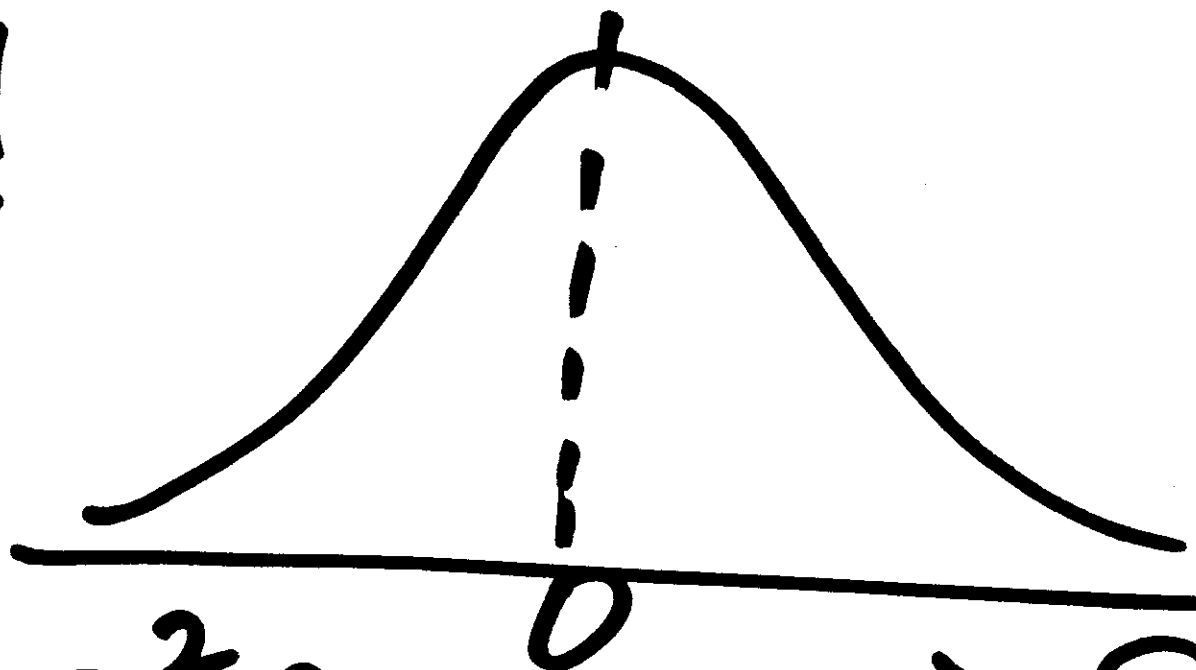
$$\ln \Omega_2 - \ln \Omega_1$$

finite

$$\frac{1}{\sqrt{2\pi}} e^{-\frac{t^2}{2}} \rightarrow -\Omega^2$$

Is this
admissible? $\rightarrow e$

No!

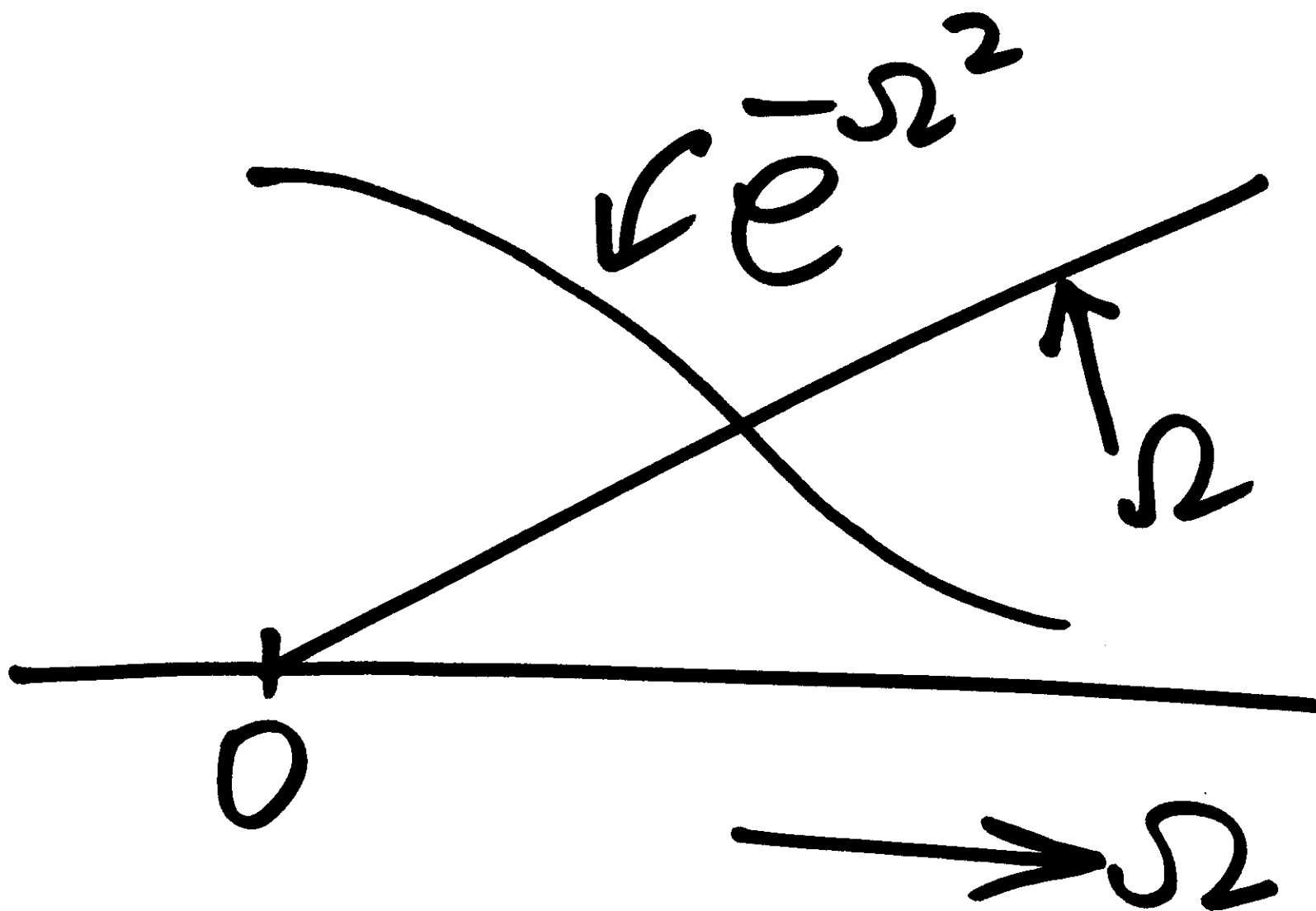


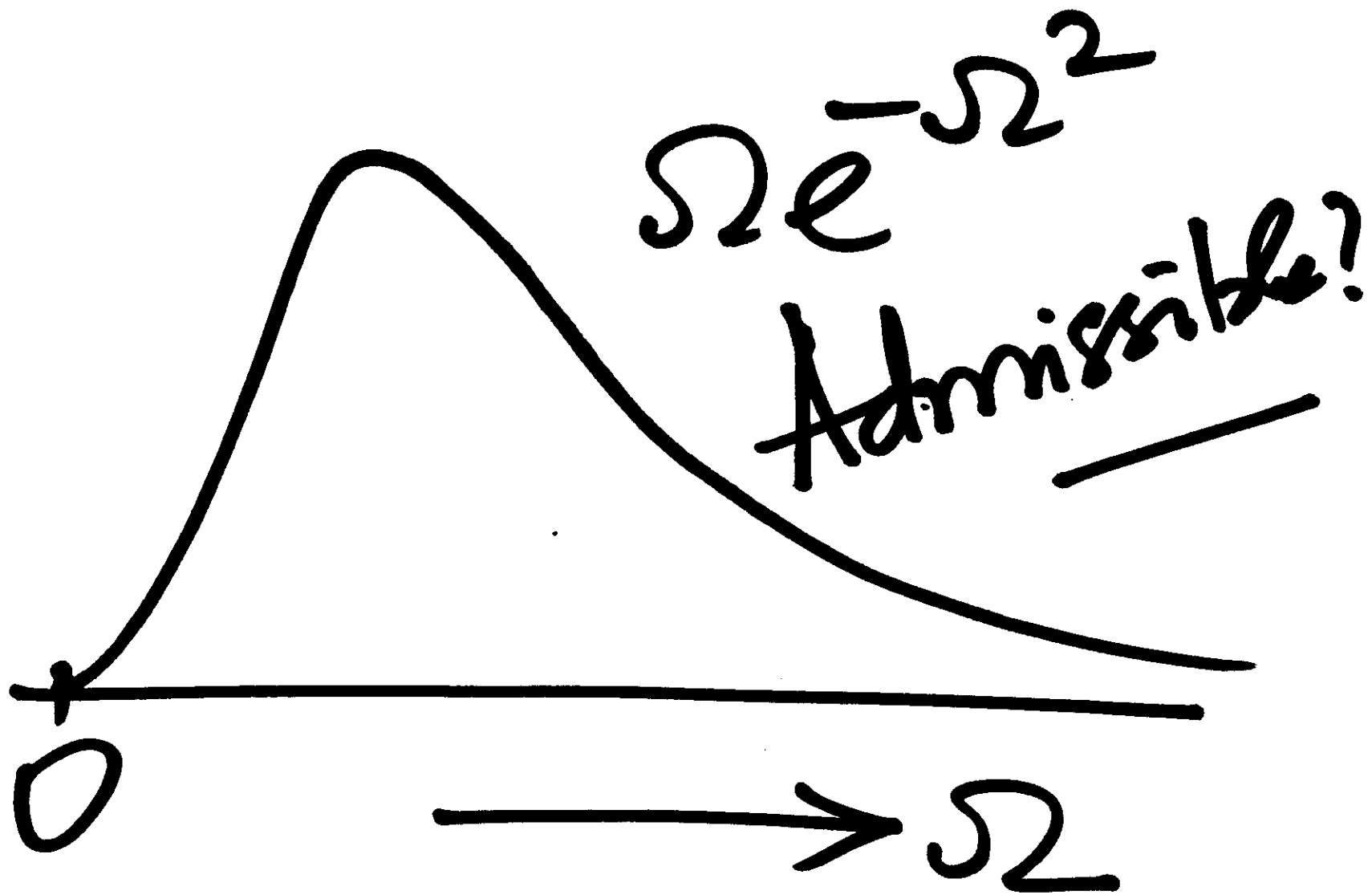
$$\int_0^1 \frac{1}{|e^{-\Omega^2}|^2} \frac{d\Omega}{\Omega} \text{ diverges}$$

Derivative of
Gaussian:

$$\frac{d}{d\Omega} e^{-\Omega^2} = -2\Omega e^{-\Omega^2}$$

$\Rightarrow \Omega e^{-\Omega^2}$





$$\int_0^{\infty} |\Omega e^{-\Omega^2}|^2 \frac{d\Omega}{\Omega}$$

$$= \int_0^{\infty} \Omega^2 e^{-2\Omega^2} \frac{d\Omega}{\Omega}$$

$$= \int_0^{\infty} e^{-2\omega^2} \cdot \omega d\omega$$

$$\omega^2 = \lambda$$
$$d\lambda = 2\omega d\omega$$

$$= \int_0^{\infty} e^{-2\lambda} \cdot \frac{1}{2} d\lambda$$

Converges

Exercises:

1. Calculate (if need
be numerically/
approx) $TBP_{\sigma_t \sigma_\Omega^2}$ of $-\Omega^2$
 Ωe

$$\frac{d}{d\Omega} e^{-\Omega^2}$$

(to within constants)

$$\frac{d}{d\Omega}$$

is equivalent
to multiply
by t

The inverse Fourier
transform of $\Omega e^{-\Omega^2}$
has the same
(exercise) form.

2. Exercise:

Consider

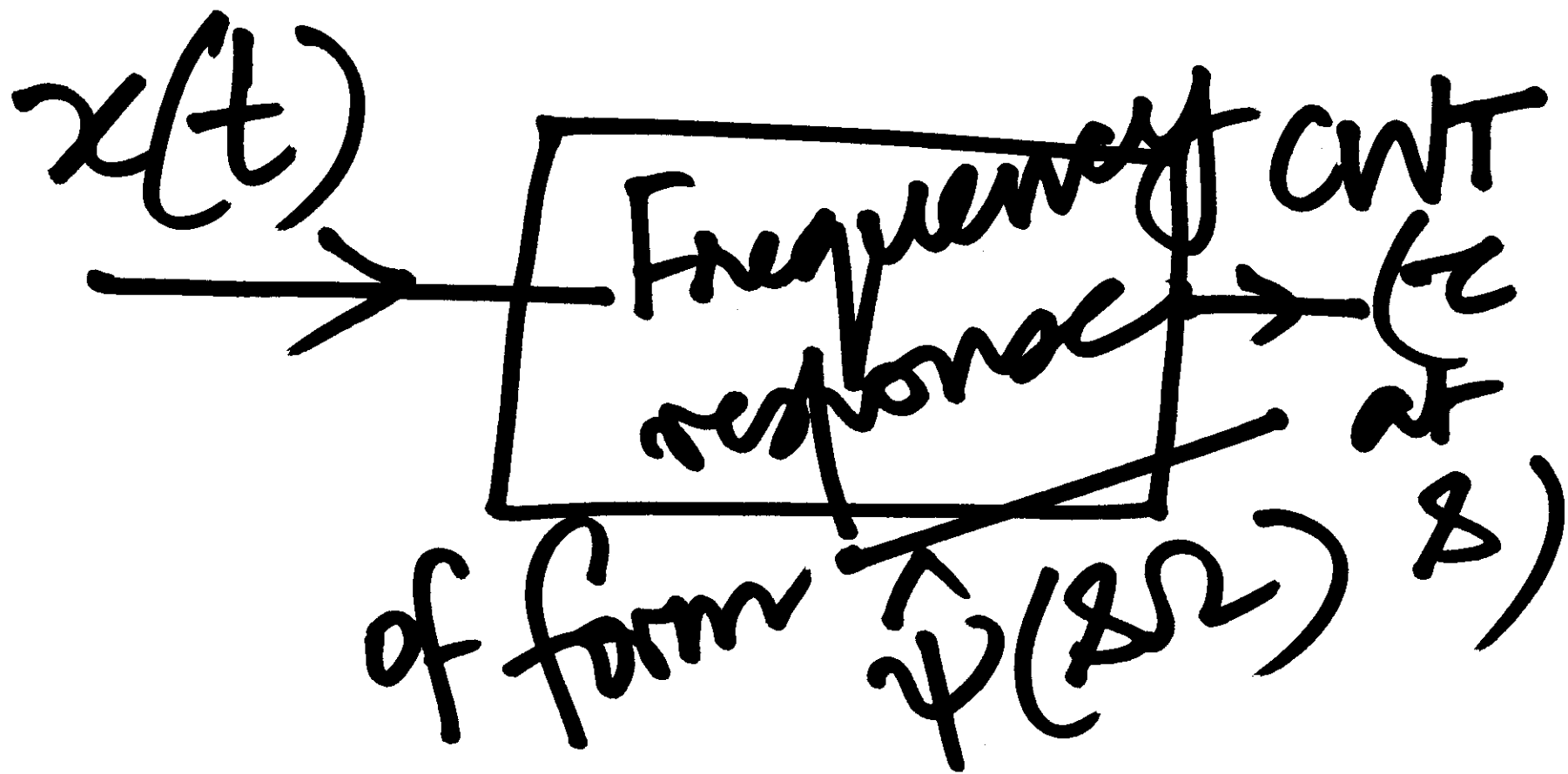
$$\frac{d^2}{d\Omega^2} e^{-\Omega^2}$$



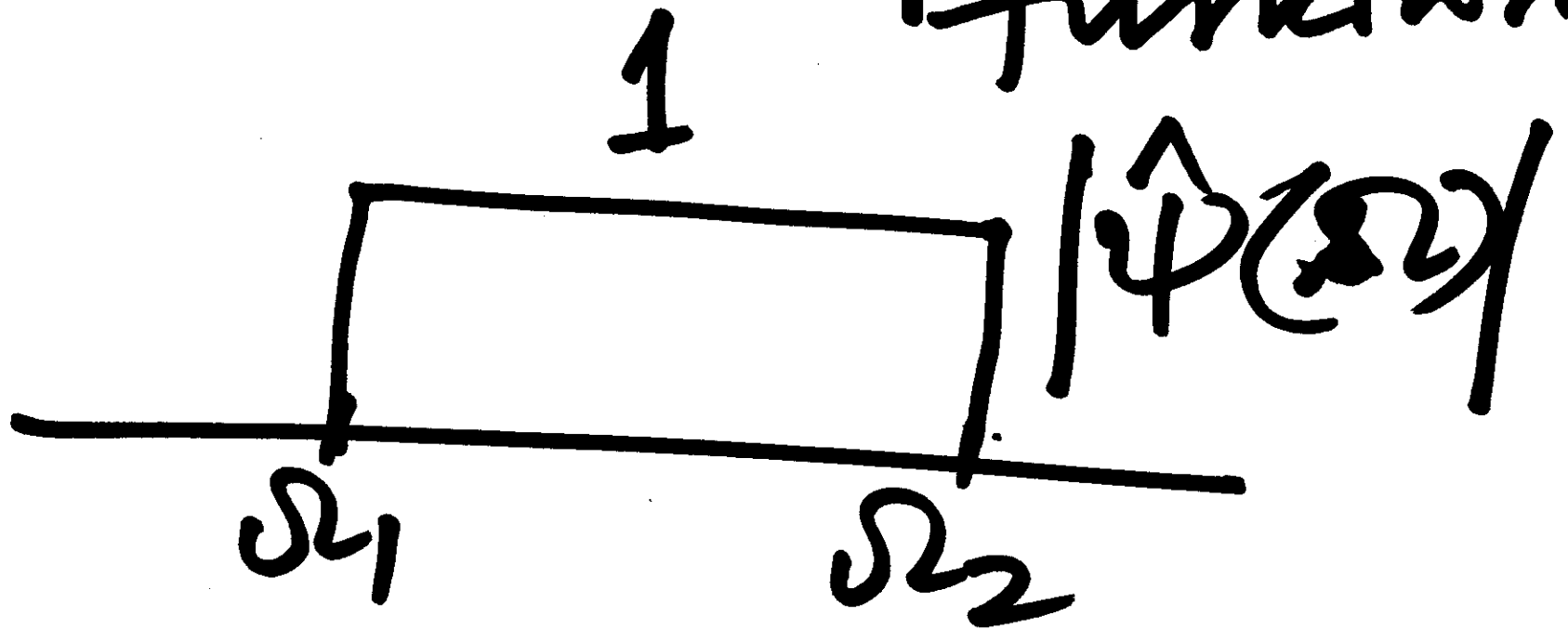
O ↓ derivative
would look like
a Mexican hat

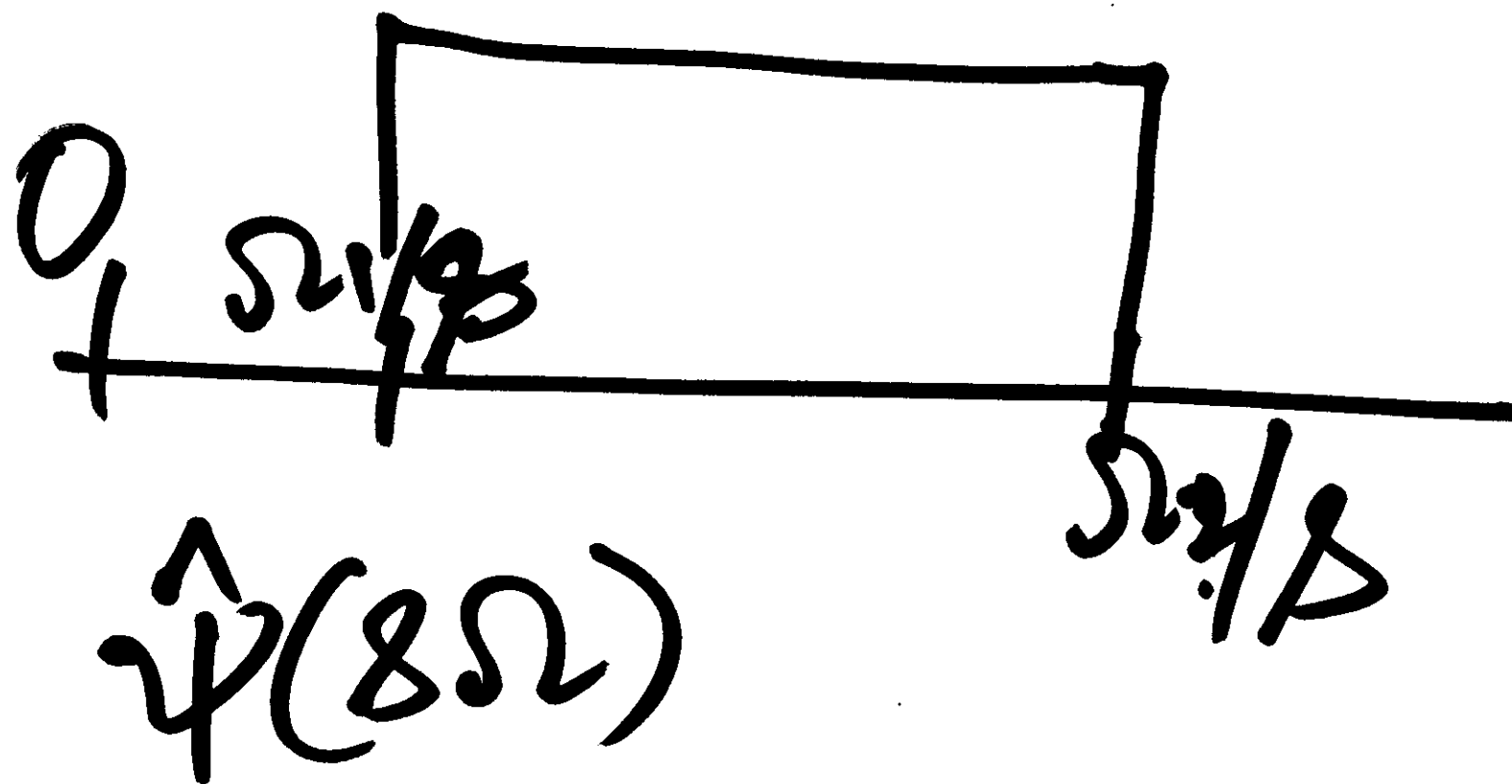
Admissibility is
essentially a
"Bandpass"
requirement

When can we
discretize $\mathcal{S}$?



Ideal bandpass
function





We could, in general
allow

$$f = a_0$$

k ; all integers;
 $a_0 \geq 1$