ME-662 CONVECTIVE HEAT AND MASS TRANSFER

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LECTURE-4 SCALAR TRANSPORT EQUATIONS

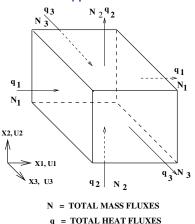
LECTURE-4 SCALAR TRANSPORT EQUATIONS

- Law of Mass Conservation for a Specie in a Mixture (Mass Transfer Equation)
- 1st Law of Thermodynamics (Energy Equation)

Mass Transfer Equation - L4($\frac{1}{17}$)

For Specie k in a Mixture Rate of accumulation of mass ($\dot{M}_{k,ac}$) = Rate of mass in ($\dot{M}_{k,in}$) - Rate of mass out ($\dot{M}_{k,out}$) + Rate of generation within CV ($R_{k,cv}$)

$$\begin{split} \dot{M}_{k,ac} &= \frac{\partial(\rho_k \Delta V)}{\partial t} \\ \dot{M}_{k,in} &= N_{1,k} \Delta A_1 \mid_{x_1} + \\ N_{2,k} \Delta A_2 \mid_{x_2} + N_{3,k} \Delta A_3 \mid_{x_3} \\ \dot{M}_{k,out} &= N_{1,k} \Delta A_1 \mid_{x_1 + \Delta x_1} \\ + N_{2,k} \Delta A_2 \mid_{x_2 + \Delta x_2} + \\ N_{3,k} \Delta A_3 \mid_{x_3 + \Delta x_3} \end{split}$$



 $\rho_k = \text{Specie Density}$ Substitute, Divide each term by ΔV and Let $\Delta x_1, \Delta x_2, \Delta x_3 \rightarrow 0$

Mass Transfer Equation - I L4($\frac{2}{17}$)

$$\frac{\partial(\rho_k)}{\partial t} + \frac{\partial(N_{1,k})}{\partial x_1} + \frac{\partial(N_{2,k})}{\partial x_2} + \frac{\partial(N_{3,k})}{\partial x_3} = R_k$$
(1)

Now, the total mass transfer flux $N_{i,k}$ in direction i is the sum of *Convective Flux* ($\rho_k u_i$) due to bulk fluid motion and *Diffusion* Flux ($m''_{i,k}$) due to density difference. Thus,

$$N_{i,k} = \rho_k \, u_i + m_{i,k}^{''}$$
 (2)

where

$$m_{i,k}^{''} = -D \frac{\partial \rho_k}{\partial x_i}$$
 (3)

where D (m^2/s) is the mass-diffusivity

Mass Transfer Equation - II L4($\frac{3}{17}$) Substituting for $N_{i,k}$ gives

$$\frac{\partial(\rho_k)}{\partial t} + \frac{\partial(\rho_k u_1)}{\partial x_1} + \frac{\partial(\rho_k u_2)}{\partial x_2} + \frac{\partial(\rho_k u_3)}{\partial x_3}$$
$$= \frac{\partial}{\partial x_1} \left(D \frac{\partial \rho_k}{\partial x_1} \right) + \frac{\partial}{\partial x_2} \left(D \frac{\partial \rho_k}{\partial x_2} \right) + \frac{\partial}{\partial x_3} \left(D \frac{\partial \rho_k}{\partial x_3} \right) + R_k \quad (4)$$

Define Mass Fraction ω_k

$$\omega_{k} = \frac{\rho_{k}}{\rho_{m}} \sum_{\substack{\text{all species}}} \omega_{k} = 1 \text{ and } \sum \rho_{k} = \rho_{m} \quad (5)$$

$$\frac{\partial(\rho_{m} \,\omega_{k})}{\partial t} + \frac{\partial(\rho_{m} \,u_{j} \,\omega_{k})}{\partial x_{j}} = \frac{\partial}{\partial x_{j}} \left(\rho_{m} \,D \,\frac{\partial\omega_{k}}{\partial x_{j}}\right) + R_{k} \quad (6)$$
Summation gives Bulk Mass Conservation because $\sum \omega_{k} = 1$
and hence, $\sum R_{k} = 0$. Also, $\sum N_{i,k} = \rho_{m} u_{i} + \sum \phi_{m} \dot{m}_{i,k}^{"} = \rho_{m} u_{i} + \sum \phi_{m} \dot{m}_{i,k}^{"} = \rho_{m} u_{i} + \rho_{m} u$

Mass Diffusivity L4($\frac{4}{17}$)

- The mass diffusivity is *defined* only for a Binary Mixture of two fluids 1 and 2 as D_{12} .
- In Multicomponent Gaseous Mixtures, however, diffusivities for pairs of species are nearly equal and a single symbol D suffices for all species.
- Incidentally, in turbulent flows, this assumption of equal (effective) diffusivities holds even greater validity as will be shown later

1st Law of Thermodynamics L4($\frac{5}{17}$)

In Rate Form (W / m^3), the 1st Law of Thermodynamics reads as

$$\dot{E} = \dot{Q}_{conv} + \dot{Q}_{cond} + \dot{Q}_{gen} - \dot{W}_s - \dot{W}_b$$
(7)

where

- \dot{E} = Rate of Change of Energy of the CV
- \dot{Q}_{conv} = Net Rate of Energy transferred by Convection
- \dot{Q}_{cond} = Net Rate of Energy transferred by Conduction
- \dot{Q}_{gen} = Net Rate of Volumetric Heat Generation within CV
 - \dot{W}_{s} = Net Rate of Work Done by Surface Forces
 - \dot{W}_b = Net Rate of Work Done by Body Forces

Each term will now be represented by a mathematical expression.

Rate of Change L4($\frac{6}{17}$)

$$\dot{E} = \frac{\partial(\rho_m \, e^o)}{\partial t} \tag{8}$$

$$e^o = e_m + \frac{V^2}{2} = h_m - \frac{p}{\rho_m} + \frac{V^2}{2} \tag{9}$$

where

 $e_m = \text{Mixture Specific Energy (J / kg)}$ $h_m = \text{Mixture Specific Enthalpy (J / kg)}$ $V^2 = u_1^2 + u_2^2 + u_3^2 (J / kg)$ $\rho_m = \sum \rho_k = \text{Mixture Density (kg / m^3)}$

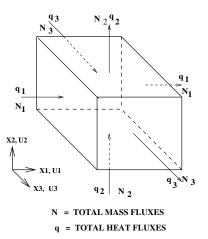
Contributions to e^0 from other forms (Potential, Electro-Magnetic etc) of energy are neglected.

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Thermodynamic Convention L4($\frac{7}{17}$)

Convention:

- Heat Flow into CV is Positive whereas Heat Flow out of CV is Negative.
- Also, all species are transported at *Mixture Velocity*



Net Convection L4($\frac{8}{17}$) Since $\sum N_{i,k} = \rho_m u_i$ $\dot{Q}_{conv} = -\frac{\partial \sum (N_{j,k} e_k^o)}{\partial x_j} = -\frac{\partial}{\partial x_j} \left[(\sum N_{j,k} h_k) + \rho_m u_j (-\frac{p}{\rho_m} + \frac{V^2}{2}) \right]$

Now, following definition of $N_{j,k}$ and noting $\sum \omega_k h_k = h_m$,

$$\sum N_{j,k} h_k = \sum (\rho_m u_j \omega_k + m_{j,k}'') h_k$$
$$= \rho_m u_j h_m + \sum m_{j,k}'' h_k$$
(11)

Hence, after some algebra

$$\dot{Q}_{conv} = -\frac{\partial(\rho_m \, u_j \, e^o)}{\partial x_j} - \frac{\partial(\sum m''_{j,k} \, h_k)}{\partial x_j}$$
(12)

Net Conduction L4($\frac{9}{17}$)

Similarly, from Fourier's Law of Conduction

$$\dot{\mathsf{Q}}_{cond} = -rac{\partial \mathbf{q}_{j}}{\partial \mathbf{x}_{j}} = rac{\partial}{\partial \mathbf{x}_{j}} \left[\mathbf{k}_{m} \, rac{\partial \mathbf{T}}{\partial \mathbf{x}_{j}}
ight]$$

where k_m = Mixture Conductivity

(13)

Volumetric Generation L4($\frac{10}{17}$)

$$\dot{Q}_{gen} = \dot{Q}_{chem} + \dot{Q}_{rad}$$
 (14)

- The Chemical Energy is positive for exothermic reactions and negative for endothermic reactions.
- 2 Evaluation of \dot{Q}_{chem} depends on the chemical reaction model employed in a particular situation.
- The Q_{rad} term represents the Net Radiation Exchange between the CV and its surroundings. Evaluation of this term, in general, requires solution of integro-differential equations
- When absorptivity (a) and scattering coefficient (s) are large,

$$\dot{\mathbf{Q}}_{rad} = \frac{\partial}{\partial \mathbf{x}_j} \left[\mathbf{k}_{rad} \ \frac{\partial T}{\partial \mathbf{x}_j} \right] \quad \mathbf{k}_{rad} = \frac{16 \sigma T^3}{a+s}$$
(15)

where σ is the Stefan-Boltzmann constant

Work-Done by Forces-I L4($\frac{11}{17}$)

Convention: The work done on the CV is negative,

$$-\dot{W}_{s} = \frac{\partial}{\partial x_{1}} [\sigma_{1} u_{1} + \tau_{12} u_{2} + \tau_{13} u_{3}]$$

$$+ \frac{\partial}{\partial x_{2}} [\tau_{21} u_{1} + \sigma_{2} u_{2} + \tau_{23} u_{3}]$$

$$+ \frac{\partial}{\partial x_{3}} [\tau_{31} u_{1} + \tau_{32} u_{2} + \sigma_{3} u_{3}]$$
(16)

$$-W_{b} = \rho_{m} \left(B_{1} u_{1} + B_{2} u_{2} + B_{3} u_{3} \right)$$
(17)

where W_s = Stress Work and W_b = Body-Force Work Further use of *differentiation of product* gives (see next slide)

Work-Done by Forces-II L4($\frac{12}{17}$)

$$-(\dot{W}_{s}+\dot{W}_{b}) = u_{1} \left[\frac{\partial\sigma_{1}}{\partial\mathbf{x}_{1}}+\frac{\partial\tau_{21}}{\partial\mathbf{x}_{2}}+\frac{\partial\tau_{31}}{\partial\mathbf{x}_{3}}+\rho_{m}B_{1}\right]$$
(18)
$$= u_{2} \left[\frac{\partial\tau_{12}}{\partial\mathbf{x}_{1}}+\frac{\partial\sigma_{2}}{\partial\mathbf{x}_{2}}+\frac{\partial\tau_{32}}{\partial\mathbf{x}_{3}}+\rho_{m}B_{2}\right]$$
(19)
$$+ u_{3} \left[\frac{\partial\tau_{13}}{\partial\mathbf{x}_{1}}+\frac{\partial\tau_{23}}{\partial\mathbf{x}_{2}}+\frac{\partial\sigma_{3}}{\partial\mathbf{x}_{3}}+\rho_{m}B_{3}\right]$$
(20)
$$+ \sigma_{1}\frac{\partialu_{1}}{\partial\mathbf{x}_{1}}+\sigma_{2}\frac{\partialu_{2}}{\partial\mathbf{x}_{2}}+\sigma_{3}\frac{\partialu_{3}}{\partial\mathbf{x}_{3}}$$
(21)
$$+ \tau_{12} \left(\frac{\partialu_{1}}{\partial\mathbf{x}_{2}}+\frac{\partialu_{2}}{\partial\mathbf{x}_{1}}\right)+\tau_{13} \left(\frac{\partialu_{1}}{\partial\mathbf{x}_{3}}+\frac{\partialu_{3}}{\partial\mathbf{x}_{1}}\right)$$

$$+ \tau_{23} \left(\frac{\partialu_{2}}{\partial\mathbf{x}_{3}}+\frac{\partialu_{3}}{\partial\mathbf{x}_{2}}\right)$$
(22)

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Complementary stress $\tau_{12} = \tau_{21}$ etc are recognised (see next slide) March 16, 2010

Work-Done by Forces-III L4($\frac{13}{17}$)

Multipliers of u_1 , u_2 , u_3 in equations 18, 19 and 20 are simply RHS of Momentum equations (See Lecture 3, slides 13-14-15). They are replaced by LHS of Momentum equations. Hence,

Equations 18, 19, 20 =
$$\rho_m \left[u_1 \frac{Du_1}{Dt} + u_2 \frac{Du_2}{Dt} + u_3 \frac{Du_3}{Dt} \right]$$

= $\rho_m \frac{D}{Dt} \left(\frac{V^2}{2} \right)$ (23)

Work-Done by Forces-IV L4($\frac{14}{17}$)

Similarly, using Stokes's Laws , equations 21, and 22 can be written as

Equations 21,
$$22 = \mu \Phi_v - \rho \bigtriangledown V$$
 (24)

where Viscous Dissipation Function is

$$\Phi_{\nu} = 2 \left[\left(\frac{\partial u_1}{\partial x_1} \right)^2 + \left(\frac{\partial u_2}{\partial x_2} \right)^2 + \left(\frac{\partial u_3}{\partial x_3} \right)^2 \right] \\ + \left(\frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right)^2 + \left(\frac{\partial u_1}{\partial x_3} + \frac{\partial u_3}{\partial x_1} \right)^2 + \left(\frac{\partial u_3}{\partial x_2} + \frac{\partial u_2}{\partial x_3} \right)^2$$
(25)

Work-Done by Forces-V L4($\frac{15}{17}$)

Hence, from equations 23 and 24, we have

$$-(\dot{W}_{s}+\dot{W}_{b})=\rho_{m}\frac{D}{Dt}\left(\frac{V^{2}}{2}\right)+\mu\Phi_{v}-\rho \bigtriangledown V \qquad (26)$$

where

$$\nabla \cdot V = \frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3}$$
 (27)

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Summary L4($\frac{16}{17}$ **)**

$$\dot{E} = \dot{Q}_{conv} + \dot{Q}_{cond} + \dot{Q}_{gen} - \dot{W}_s - \dot{W}_b$$
(28)

where

$$\dot{E} = \frac{\partial(\rho_m e^o)}{\partial t}$$
(29)

$$\dot{Q}_{conv} = -\frac{\partial(\rho_m u_j e^o)}{\partial x_j} - \frac{\partial(\sum m''_{j,k} h_k)}{\partial x_j}$$
(30)

$$\dot{Q}_{cond} = \frac{\partial}{\partial x_j} \left[k_m \frac{\partial T}{\partial x_j} \right]$$
(31)

$$\dot{Q}_{gen} = \dot{Q}_{chem} + \dot{Q}_{rad}$$
(32)

$$-(\dot{W}_s + \dot{W}_b) = \rho_m \frac{D}{Dt} \left(\frac{V^2}{2} \right) + \mu \Phi_v - p \nabla \cdot V$$
(33)

Final Energy Equation L4($\frac{17}{17}$)

$$\frac{\partial(\rho_m \, \mathbf{e}^{\mathbf{o}})}{\partial t} + \frac{\partial(\rho_m \, u_j \, \mathbf{e}^{\mathbf{o}})}{\partial \mathbf{x}_j} = \frac{\partial}{\partial \mathbf{x}_j} \left[\mathbf{k}_m \, \frac{\partial \mathbf{T}}{\partial \mathbf{x}_j} \right] - \frac{\partial(\sum m_{j,k}^{\prime\prime} \, \mathbf{h}_k)}{\partial \mathbf{x}_j} \\ + \frac{D}{D \, t} \left[\frac{\mathbf{V}^2}{2} \right] - \mathbf{p} \, \bigtriangledown \, \mathbf{V} + \mu \, \Phi_v \\ + \dot{\mathbf{Q}}_{chem} + \dot{\mathbf{Q}}_{rad}$$
(34)

or, using definition of e^o,

$$\rho_{m} \frac{D h}{D t} = \frac{\partial}{\partial x_{j}} \left[k_{m} \frac{\partial T}{\partial x_{j}} \right] - \frac{\partial (\sum m_{j,k}^{"} h_{k})}{\partial x_{j}} + \mu \Phi_{v} + \frac{D p}{D t} + \dot{Q}_{chem} + \dot{Q}_{rad}$$
(35)

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