

ME-662 CONVECTIVE HEAT AND MASS TRANSFER

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LECTURE-31 CONVECTIVE MASS TRANSFER

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- 4 Boundary Layer Forms
- 5 Simplified Forms
 - 1 Stefan Flow Model
 - 2 Couette Flow Model
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Importance of CHM L31($\frac{1}{14}$)

- 1 In several applications such as **evaporative or transpiration cooling, combustion, ablation, condensation, drying etc** , convective mass transfer assumes importance.
- 2 CHM takes place **between two dissimilar phases** such as liquid - gas or solid - gas. Hence, CHM is also referred to as **Inter-phase Mass Transfer**
- 3 In the absence of fluid motion, mass transfer takes place by **diffusion** only. Analogous to **heat conduction** , diffusion mass transfer takes place **within a single phase** ; gas, liquid or solid
- 4 Examples of diffusion mass transfer are: **Leakage of hydrogen** by diffusion through the solid container wall, Penetration of carbon through solid iron during **hardening** , diffusion of water vapour through solid during **drying** , diffusion of **sugar through stagnant liquid tea in a cup**

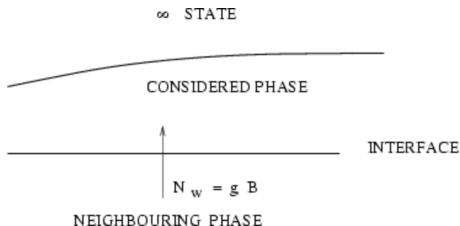
Main Task - L31($\frac{2}{14}$)

- 1 CHM takes place due to concentration gradients of the transferred species
- 2 The main task is to calculate Interface mass transfer flux N_w ($\text{kg}/\text{m}^2\text{-s}$) from

$$N_w = g B$$

where g ($\text{kg}/\text{m}^2\text{-s}$) is mass transfer coefficient and B is dimensionless driving force

- 3 N_w and g have same units



N_w is positive when mass transfer takes place from the *neighbouring phase* into the *considered phase* across the interface & *vice versa*

Concentration defined - L31($\frac{3}{14}$)

- ① Both neighboring and considered phases are mixtures comprising several species k . The proportion of k in the mixture can be defined in several ways.

- ② Mass Fraction

$$\omega_k \equiv \frac{m_k}{m_{mix}} = \frac{\rho_k}{\rho_{mix}}$$

- ③ Mole Fraction

$$x_k \equiv \frac{n_k}{n_{mix}} = \frac{p_k}{p_{mix}}$$

where n = number of moles and p_k is partial pressure.

- ④ Concentration

$$[k] \equiv \frac{\rho_k}{M_k} = \frac{p_k}{R_u T} = \frac{1}{V_k} \left(\frac{\text{kmol}}{\text{m}^3} \right)$$

where M_k is molecular weight.

- ⑤ Most often, we shall prefer mass fraction ω_k which represents proportion by mass.

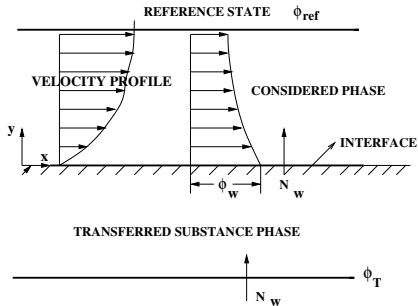
Driving force defined - L31($\frac{4}{14}$)

Unlike in heat transfer, in mass transfer, 3 states must be considered:

- 1 Reference state (ref) far into the *Considered Phase*
- 2 Interface state (w)
- 3 Trans Subs state (T) deep into *Neighbouring Phase*

Dimensionless Driving force B is

$$B = \frac{\phi_{ref} - \phi_w}{\phi_w - \phi_T}$$



where ϕ is a **Conserved property**. ϕ can be formed from ω_k , x_k or from $[k]$.

Types of Mass Transfer - L31($\frac{5}{14}$)

In general, there are 3 types of mass transfer

- 1 *Mass Transfer without heat transfer* (no chemical reaction)
- 2 *Mass Transfer with heat transfer* (no chemical reaction)
- 3 *Mass Transfer with heat transfer and chemical reaction*

In each case, Conserved Property Φ must be appropriately defined. *Any eqn of the form*

$$\frac{\partial(\rho_m \Phi)}{\partial t} + \frac{\partial}{\partial x_j} \left[\rho_m u_j \Phi - \Gamma_\Phi \frac{\partial \Phi}{\partial x_j} \right] = S_\Phi$$

in which $S_\Phi = 0$, Φ is called a conserved property.

Transport Eqns - L31($\frac{6}{14}$)

From lectures 2 and 3

$$\frac{\partial(\rho_m)}{\partial t} + \frac{\partial(\rho_m u_j)}{\partial x_j} = 0 \quad (\text{Bulk mass})$$

$$\rho_m \frac{D u_j}{D t} = - \frac{\partial p}{\partial x_j} + \frac{\partial}{\partial x_j} \left[\mu_{m,eff} \frac{\partial u_j}{\partial x_j} \right] + \rho_m B_j + \frac{\partial}{\partial x_j} \left[\mu_{m,eff} \frac{\partial u_j}{\partial x_j} \right] \quad (\text{Momentum})$$

$$\rho_m \frac{D \omega_k}{D t} = \frac{\partial}{\partial x_j} \left(\rho_m D_{eff} \frac{\partial \omega_k}{\partial x_j} \right) + R_k \quad (\text{Species Transfer})$$

$$\rho_m \frac{D h_m}{D t} = \frac{\partial}{\partial x_j} \left[k_{m,eff} \frac{\partial T}{\partial x_j} \right] - \frac{\partial(\sum m''_{j,k} h_k)}{\partial x_j} + \frac{D p}{D t} + \dot{Q}_{chem} + \dot{Q}_{rad} + \mu_{eff} \Phi_v \quad (\text{Energy})$$

Boundary Layer Model - L31($\frac{7}{14}$)

From **Near-interface** considerations

$$\frac{\partial(\rho_m \Psi)}{\partial t} + \frac{\partial(\rho_m u \Psi)}{\partial x} + \frac{\partial(\rho_m v \Psi)}{\partial y} = \frac{\partial}{\partial y} \left[\Gamma_\Psi \frac{\partial \Psi}{\partial y} \right] + S_\Psi$$

Ψ	Γ_Ψ	S_Ψ
1	0	0
u	$\mu_{m,eff}$	$- dp / dx + B_x$
ω_k	$\rho_m D_{eff}$	R_k
h_m	$k_{m,eff} / c\rho_m$	$Dp/Dt + \dot{Q}_{rad} + \dot{Q}_{others}$ $-\partial(\sum m''_{y,k} h_k) / \partial y + \mu_{eff} (\partial u / \partial y)^2$

where $m''_{y,k} = -\rho_m D_{eff} \partial \omega_k / \partial y$.

Comments on BL model - L31($\frac{8}{14}$)

- ① In the above eqns, when chemical reaction is present, species enthalpy is defined as

$$h_k(T) \equiv h_{f,k}^0 + \Delta h_{s,k} = h_{f,k}^0 + \int_{T_{ref}}^T c_{p,k} dT$$

where $h_{f,k}^0$ is **Enthalpy of formation** .

- ② Also, from Fick's law of mass diffusion

$$m''_{y,k} = - D_{eff} \frac{\partial \rho_k}{\partial y} = - \rho_m D_{eff} \frac{\partial \omega_k}{\partial y}$$

- ③ BL model is **an idealisation of the transport eqns** . It still involves simultaneous soln of several coupled differential eqns which requires computer solutions.
- ④ To avoid this, we shall postulate simpler models to serve a limited purpose but **the solutions from model eqns can reveal the tendencies inherent in the complete BL model.**

Element Transport Eqn - 1 - L31($\frac{9}{14}$)

- 1 Chemical reactions obey the *element conservation principle*. Thus, consider a mixture comprising CH_4 , O_2 , H_2 , H_2O , CO_2 , CO , N_2 , NO , and O .
- 2 Then, the mass fraction (symbol η) of elements C, H, O, and N in the mixture will be related to ω_k as

$$\eta_C = \frac{12}{16} \omega_{CH_4} + \frac{12}{44} \omega_{CO_2} + \frac{12}{28} \omega_{CO},$$

$$\eta_H = \frac{4}{16} \omega_{CH_4} + \frac{2}{2} \omega_{H_2} + \frac{2}{18} \omega_{H_2O},$$

$$\eta_O = \frac{32}{44} \omega_{CO_2} + \frac{16}{28} \omega_{CO} + \frac{32}{32} \omega_{O_2} + \frac{16}{18} \omega_{H_2O}$$

$$+ \frac{16}{16} \omega_O + \frac{16}{30} \omega_{NO} \quad \text{and} \quad \eta_N = \frac{14}{30} \omega_{NO} + \frac{28}{28} \omega_{N_2}$$

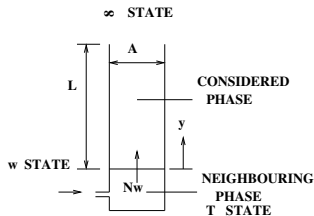
Element Transport Eqn - 2 - L31($\frac{10}{14}$)

- 1 In general, therefore, the *mass fraction* η_α of element α is $\eta_\alpha = \sum_k \eta_{\alpha,k} \omega_k$ where $\eta_{\alpha,k} = M_\alpha / M_k$ is the mass fraction of element α in the species k .
- 2 Just as the species are convected, diffused, and generated or destroyed, the elements can also be considered to have been convected and diffused, but *they can never be destroyed or generated* because of the principle of element conservation.
- 3 Thus, the transport equation for any element α will have no source term.

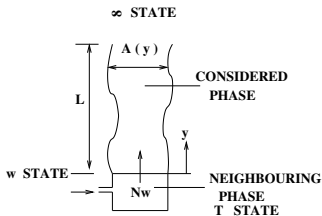
$$\frac{\partial(\rho_m \eta_\alpha)}{\partial t} + \frac{\partial(\rho_m u_j \eta_\alpha)}{\partial x_j} = \frac{\partial}{\partial x_j} (\rho_m D \frac{\partial \eta_\alpha}{\partial x_j})$$

where it is assumed that D for the elements is same as that for the species. η_α is always a conserved property.

Stefan Flow Model - L31($\frac{11}{14}$)



(a) CONSTANT AREA MODEL



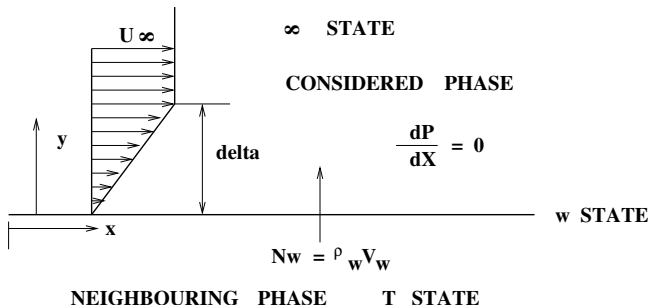
(b) VARIABLE AREA MODEL

In this 1D model, we set $u = dp/dx = 0$ and $A = A(y)$. Then, allowing for area change in y -direction, the BL model Eqn will transform to

$$\frac{d}{dy} [N_{\psi,y} A] = \frac{d}{dy} \left[\rho_m v A \psi - \Gamma_{\psi} A \frac{d\psi}{dy} \right] = A S_{\psi}$$

where $\mu_{m,t} = D_t = k_{m,t} = 0$ **Laminar diffusion only**.

Couette Flow Model - L31($\frac{12}{14}$)



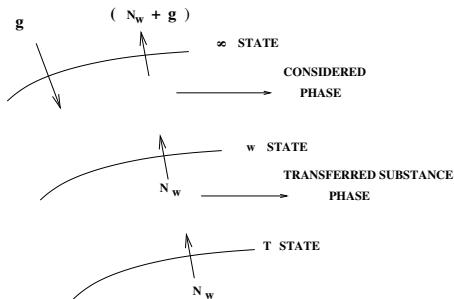
In this 1D model, we again set $dp/dx = 0$, but $u = C \times y$, $\partial/\partial x = 0$ and $A = \text{const}$. Then, in this idealisation, BL model reads as

$$\frac{d}{dy} [N_{\psi,y}] = \frac{d}{dy} \left[\rho_m v \psi - \Gamma_{\psi} \frac{d\psi}{dy} \right] = S_{\psi}$$

This model permits study of effect of turbulent $\Gamma_{\psi,t}$ and of property variations

Reynolds Flow Model - L31($\frac{13}{14}$)

- 1 This is an Algebraic model^a - no differential eqn
- 2 The model postulates inward flux g at imaginary surface in the ∞ state carrying with it properties of ∞ -state and an Outward flux $(N_w + g)$, at the same imaginary surface, carrying with it properties of the w -state



Thus, there is no net mass creation between w and ∞ states. The model claims to account for all the effects produced at the w surface in a real boundary layer model.

^aSpalding D B Convective Mass Transfer, Edward Arnold Ltd, London (1963)

Comment on the models - L31($\frac{14}{14}$)

- 1 Like in heat transfer, the greater part of the resistance to mass transfer is confined to near-wall region.
- 2 Hence, boundary layer flow model should suffice. But, this requires computer solutions of the governing eqns.
- 3 To avoid this, and to establish validity of the $N_w = g \times B$ simplified models are invoked.
- 4 Stefan flow and Couette flow models are One-dimensional and hence, analytically derived *closed form* solutions are possible.
- 5 Reynolds flow model is the simplest, being algebraic.
- 6 In the lectures to follow, the model solutions will be developed for the 3-types of mass transfer